Analytical Approaches in Investigating the Kinetics of Water-Molecule Complexes in Tropospheric Reactions

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of Water-Molecule Complexes in
Tropospheric Reactions

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A thesis submitted to the faculty of
Brigham Young University
in partial fulfillment of the requirements for the degree of
Master of Science

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ABSTRACT

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Ozone is a heavily monitored pollutant. Ozone is not directly emitted into the atmosphere, but rather the product of chemical reactions. Ground level ozone occurs when nitrogen oxides (NOₓ) and volatile organic compounds (VOCs) react with each other in the presence of sunlight. The primary precursors of ozone are anthropogenically emitted, and as a result, tropospheric ozone has cost millions of dollars in damages and has hurt the health of countless people.

This thesis is a collection of work that aims to provide insight into atmospheric reactions that result in tropospheric ozone and the instrumentation to study such reactions. While these reactions are well studied, this research is novel in its attempt to understand water vapor’s influence in tropospheric ozone reactions. As the troposphere continues to get warmer and wetter from global climate change, water vapor will play a larger role in tropospheric reactions, which in turn may perturb the global reactions.

Work is presented on the self-reaction of β-hydroxyethyl peroxy radical (β-HEP), an ozone precursor, and the increase in reaction rate catalyzed by water vapor. β-HEP serves as a model system for understanding the roles of water vapor in perturbing the kinetics and product branching ratio of ozone forming reactions. The self-reaction rate coefficient of β-HEP was investigated between 274-296 K with 1.0 × 10¹⁵ to 2.5 × 10¹⁷ molecules cm⁻³ of water vapor at 200 Torr total pressure by slow-flow laser flash photolysis coupled with UV time-resolved spectroscopy and long-path, wavelength-modulated, diode-laser spectroscopy. The overall rate constant is expressed as the product of temperature-dependent and water vapor-dependent terms giving k(T,H₂O) = 7.8 × 10⁻¹⁴ e^{8.2±2.5 kJ/RT} (1 + 1.4 × 10⁻³⁴ × e^{92±11 kJ/RT} [H₂O]). The results suggest that formation of a β-HEP-H₂O complex is responsible for the observed water vapor enhancement of the self-reaction rate coefficient.

A new discharge flow mass-spectrometer was engineered in collaboration with the California Institute of Technology and NASA’s Jet Propulsion Laboratory. This instrument allows for rapid study of water vapor influence on the kinetics of atmospheric reactions. This instrument will be used in further studying the β-HEP + NO reaction as a function of water vapor concentration.

Keywords: discharge flow mass-spectrometer, atmospheric kinetics, ozone, β-hydroxyethyl peroxy
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1 INTRODUCTION

1.1 The Atmosphere and Us

In the cold London December of 1952, a series of events lead to the death of almost 12,000 people. No shots were fired, no knives were brandished, and in the end the culprit was themselves. The Great London Smog of 1952 came to be known as the worst air-pollution event in the United Kingdom. The cold air and windless conditions allowed fog to condense near the ground. The fog then mixed with domestic smoke, factory emissions, and vehicular exhaust. Normally the fog and pollutants would rise and be dispersed in the upper atmosphere, but the cold temperatures kept the fog near the ground. Over a four day period, the trapped airborne pollutants killed thousands of people.

At the time, pollution was not a new concept; however, the Great Smog showed that pollution can be deadly. This event impacted science, public perception, and government regulations in regards to pollution and the atmosphere. Since then, scientists have learned much more about the causes and impacts of atmospheric pollution. A pioneer in photochemistry, Philip Leighton, stated:

“One of the most striking developments of the past decade, both in the field of air pollution and in that of atmospheric chemistry, has been the recognition that photochemical reactions, produced by sunlight, may convert relatively innocuous
pollutants into substances which constitute a nuisance, create a possible health hazard, and cause economic loss to man”

The chemistry of the atmosphere is complex and varies with location, altitude, season, geology, and human activity. The ability to solve or prevent pollution is dependent upon the understanding of these atmospheric reactions. Furthermore, some solutions are not universal. Solutions to pollution problems in Utah may not apply to problems in Los Angeles, and vice versa. There needs to be an understanding of local atmospheric chemistry, along with the measurement of local pollution sources. To integrate strategies for air pollution effectively, policy makers need extensive information about key pollutants and their interactions in the atmosphere.

1.2 National Ambient Air Quality Standards

In 1963, the United States Congress voted in the Clean Air Act, which commissioned research into tropospheric pollution. In 1970, this act was further amended to consolidate all federal research, monitoring, enforcement, and standard-setting into one agency, the Environmental Protection Agency (EPA). Since then, the EPA has set the National Ambient Air Quality Standards (NAAQS) for six pollutants. These pollutants include carbon monoxide, lead, nitrogen dioxide, ozone, particulate matter, and sulfur dioxide. Standards are set for these pollutants, both as primary and secondary emissions. Primary emissions affect public health, including asthmatics, children, and the elderly. Secondary emissions also disrupt public health, along with plants, animals, buildings, and visibility. States and areas that fail to maintain the NAAQS are declared, “nonattainment areas,” and are required by law to develop plans to meet
those standards or lose government funding. These plans to attain or maintain NAAQS are known as State Implementation Plans (SIPs). In order to develop realistic SIPs, policy makers must understand the chemical formation and loss mechanisms of the six pollutants. This thesis aims to contribute to both the kinetic modeling and sampling of ozone; one of the criteria pollutants.

1.3 Ozone

The discovery of ozone dates back over 150 years. Ozone was discovered in 1840 when Schonbein noticed a distinct smell during the electrolysis of water, and named the substance ozone after the Greek word "ozein" for smell. Previous scientists had noted a similar smell after lightning struck or the smell between two arcing electrodes. These smells were thought to be the result of electricity and not a chemical reaction. It wasn’t until Leighton’s work in the mid-1900s that ozone was attributed to the ozone layer and ozone pollution in the troposphere.

Ozone is a unique molecule. In the stratosphere, ozone protects living organisms from deep UV-light ($\lambda < 300$nm). But, in the lower troposphere, ozone is a harmful pollutant. Ozone is a strong oxidizer, harming both human health and vegetation. Inhalation of ozone temporarily decreases lung capacity, inflames the airways, and makes people more susceptible to allergens. The EPA characterizes the symptoms of ozone inhalation as coughing, respiratory irritation while breathing deeply, and shortness of breath. Children exposed to ozone, have an increase of respiratory problems not only as kids, but as adults. One recent study concluded that ozone exposure has greater medium- to long-term effects on subjects than short-term affects. Ozone may also harm the growth of lungs in children. In a study of 95 urban communities in the
United States over 13 years, mortality rates increased with increased ozone levels, particularly cardiovascular and respiratory mortality.

Ozone also harms vegetation and decreases crop yield. Ozone slows photosynthesis and causes physical damage to the plant tissue. Forest areas near large cities, such as Los Angeles and Mexico City, show growth decline and visible damage from elevated ozone levels. The physical damage changes the appearance of plants, decreasing their economic value, especially of horticultural crops. In 2000, 79–121 million metric tons of crops were globally lost to ozone, a $11 – 18 billion loss to the economy. In was estimated that ozone reduced the yield of the soybean crop by 8.5–14%, the wheat crop by 3.9–15%, and the maize crop by 2.2–5.5%. In another study, Avnery et al. predict that crop yields will continue to decrease due to ozone damage, with a projected annual global loss of $12 – 13 billion in 2030.

1.4 Peroxy Radicals and the HOx and NOx Cycles

Tropospheric ozone is the byproduct of the HOx and NOx cycles. Ozone episodes were first noticed in the 1940s in Los Angeles. The pollution of Los Angeles was strongly oxidizing, eye-watering, and plant killing. These pollution episodes occurred mostly at midday with warm temperatures and bright sunshine. The 1950s produced significant achievements in understanding the role of radical chemistry on ozone production. The importance of peroxy radical (RO2) chemistry was established in the late 1970s, and forms the basis of the HOx and NOx cycles. The formation of peroxy radicals by either of these pathways sets in motion the HOx and NOx cycles.
Although N₂, O₂, and Ar make up 99.96% of the Earth's atmosphere, the concentrations of trace species drive the pollution mechanisms in the atmosphere. The lower atmosphere can be thought of as a low temperature combustion system initiated through photochemistry.¹⁹ The photolysis of trace species initiates radical propagation, terminating in closed-shell species that pollute the atmosphere. Free radicals have extremely high reactivity and accordingly have short lifetimes and, therefore, are found in the troposphere at concentrations typically less than 100 parts per trillion. Their high reactivity causes the chemistry of the troposphere to be complex, but can be simplified by discovering common reactions for the removal of organics found in the atmosphere. Two prominent cycles that are commonly used to model the troposphere are the HOₓ cycle (the reactions governing the radicals OH, HO₂, and hydrogen atoms) and the NOₓ cycle (the reactions describing the interconversion of NO and NO₂). These cycles are interconnected but discussed separately here.

*Figure 1.1 - The HOₓ cycle (Author original)*
The HO\textsubscript{x} cycle is depicted schematically in Figure 1.1. The initiator of this cycle is the hydroxyl radical (OH) initially formed by photolysis of ozone in the presence of water.\textsuperscript{20}

\[
\text{O}_3 + h\nu (\lambda < 320\text{nm}) \rightarrow \text{O}(^1\text{D}) + \text{O}_2(^1\Delta g) \quad (1.1)
\]

\[
\text{O}(^1\text{D}) + \text{H}_2\text{O} \rightarrow \text{OH} + \text{OH} \quad (1.2)
\]

OH concentrations typically range from less than 1 × 10\textsuperscript{4} molecules cm\textsuperscript{-3} at night to approx. 3 × 10\textsuperscript{6} molecules cm\textsuperscript{-3} at midday.\textsuperscript{21} OH predominately reacts with CO to form hydrogen atoms that add to O\textsubscript{2} to form HO\textsubscript{2}.

\[
\text{OH} + \text{CO} \rightarrow \text{CO}_2 + \text{H} \quad (1.3)
\]

\[
\text{H} + \text{O}_2 \rightarrow \text{HO}_2 \quad (1.4)
\]

Organic species, whether from biogenic or anthropogenic sources, are removed from the troposphere by reaction with OH, creating an alkyl radical (R).

\[
\text{OH} + \text{RH (Organics)} \rightarrow \text{H}_2\text{O} + \text{R} \quad (1.5)
\]

With 21% of the atmosphere being O\textsubscript{2}, the only significant reaction of alkyl radicals is the addition of O\textsubscript{2} to form peroxy radicals.
RO₂ and the peroxo radical, HO₂, are found in concentrations from 0.2 - 2.5 × 10⁸ molecules cm⁻³.²⁰ These molecules react with themselves and with each other, especially in pristine environments (low NOₓ). The self-reaction of one of the most prominent peroxo radicals, HOCH₂CH₂O₂, is the focus of the research in Chapter 2. In more polluted conditions (high NOₓ) the self-reactions of peroxo radicals compete with the reaction with NO to oxidize NO to NO₂ and produce an alkoxy radical (RO).

\[
\text{RO}_2 + \text{NO} \rightarrow \text{RO} + \text{NO}_2 \quad (1.7)
\]

Reaction 1.7 has a second minor pathway that produces an organic nitrate (RONO₂) and is the focus of the research described in chapter 3. Alkoxy radicals react with O₂ to form HO₂ and a carbonyl containing compound.

\[
\text{RO} + \text{O}_2 \rightarrow \text{HO}_2 + \text{Carbonyl Compound} \quad (1.8)
\]

The carbonyl compound formed will depend on the R group of the original organic molecule, but often form aldehydes. Many aldehydes are known to photolyze in the daytime to produce many other radical-chain pathways. Other carbonyl compounds are oxidized by OH as the HOₓ cycle continues. The cycle completes back to where it started when HO₂ is reduced by oxidizing NO to NO₂.
\[ \text{HO}_2 + \text{NO} \rightarrow \text{OH} + \text{NO}_2 \]  
(1.9)

Chain termination occurs with peroxide formation from the self-reaction of \text{HO}_2 and in many cross reactions of \text{RO}_2.

\[ \text{NO}_2 + h\nu (\lambda < 410\text{nm}) \rightarrow \text{NO} + \text{O}(^{3}\text{P}) \]  
(1.10)

The \textit{NO}_x cycle is depicted as a schematic in Figure 1.2. This cycle is initiated by the emission of NO into the troposphere. NO and NO\textsubscript{2} are formed when N\textsubscript{2} and O\textsubscript{2} are present together at elevated temperature (T > 2800 °F), which commonly occurs in internal combustion engines and power plant boilers. As previously mentioned, Reactions 1.7 and 1.9 oxidize NO to NO\textsubscript{2}. Additionally, ozone may oxidize NO to NO\textsubscript{2}, but at slower rates. NO\textsubscript{2} is readily photolyzed to form a radical oxygen that subsequently forms ozone.\textsuperscript{22}
\[ \text{O}^{(3P)} + \text{O}_2 + \text{M} \rightarrow \text{O}_3 + \text{M} \quad (1.11) \]

This reaction is the predominant mechanism for tropospheric ozone formation.\textsuperscript{20} For every carbon in a hydrocarbon that passes through the HO\textsubscript{x} cycle, two molecules of ozone are formed. As the sun sets and the supply of OH is depleted, these cycles do not terminate because of nighttime chemistry. NO\textsubscript{2} is no longer photolyzed at night but can react with ozone to form the nitrate radical (NO\textsubscript{3}).

\[ \text{NO}_2 + \text{O}_3 \rightarrow \text{NO}_3 + \text{O}_2 \quad (1.12) \]

Nitrate subsequently takes the role of OH in the HO\textsubscript{x} cycle in oxidizing hydrocarbons through hydrogen abstraction to alkyl radicals. Nitrate chemistry ceases at dawn because of its large absorption cross section of visible light and high photolysis yield.\textsuperscript{23}

1.5 Water Vapor

Many of these tropospheric reactions that lead to ozone formation have been thoroughly studied; however, these reactions have not been investigated in the presence of water vapor. Water molecules have the ability to complex with atmospheric molecules via hydrogen bonds. The H\textsubscript{2}O molecule then acts as a third body, and possibly alters the kinetics and branching ratios of tropospheric ozone forming reactions.
1.6 Layout of Thesis

Chapter 2 is a published work, while chapters 3 and 4 are supplemental work that is currently being prepared for publication. Chapter 2 details the work done on the kinetics of the self-reaction of \( \beta \)-hydroxyethyl peroxy radicals in the presence of water vapor. Chapter 3 constitutes the principal focus of this thesis, and outlines the engineering of an instrument capable of measuring the kinetics and product branching ratio of key atmospheric reactions. Chapter 3 outlines future applications of the instrument discussed in Chapter 3, specifically the application in studying the \( \beta \)-hydroxyethyl peroxy radical + NO reaction in the presence of water vapor.
2 WATER VAPOR ENHANCEMENT OF RATES OF PEROXY RADICAL REACTIONS

2.1 Disclaimer

_The work presented in this chapter has been accepted by the International Journal of Chemical Kinetics. It is be presented in its entirety._

2.2 Abstract

Peroxy radicals can complex with water vapor. These complexes affect tropospheric chemistry. In this study, \( \beta \)-HEP (hydroxyethyl peroxy radical) serves as a model system for investigating the effect of water vapor on the kinetics and product branching ratio of the self-reaction of peroxy radicals. The self-reaction rate coefficient was determined at 274–296 K with water vapor between 1.0\( \times \)10\(^{15}\) and 2.5\( \times \)10\(^{17}\) molecules cm\(^{-3}\) at 200 torr total pressure by slow-flow laser flash photolysis coupled with UV time-resolved spectroscopy and long-path, wavelength modulated, diode-laser spectroscopy. The overall self-reaction rate constant expressed as the product of both a temperature dependent and water vapor dependent term is

\[
k_o = 7.8 \times 10^{-14} \exp((8.3 \pm 2.5 \text{ kJ/mol})/RT) + ((13.2 \pm 1.56) \times 10^{-44} \times \exp((79.3 \pm 17.18 \text{ kJ/mol})/RT) \times [H_2O])
\]

suggesting formation of a \( \beta \)-HEP-H\(_2\)O complex is responsible for the increase in the self-reaction rate coefficient with increasing water concentration. Complex formation is supported by computational results identifying three local energy minima for the \( \beta \)-HEP-H\(_2\)O complex. As the troposphere continues to get warmer and wetter, more of the peroxy radicals present will be complexed with water. Investigating the effect of water vapor on kinetics
of atmospherically relevant radicals and determining the effects of these altered kinetics on tropospheric ozone concentrations is thus important.

2.3 Introduction

A direct relationship has been established between surface temperature and tropospheric water vapor. As the Earth continues to warm and tropospheric water vapor increases, water mediated chemistry has the potential to perturb the chemistry of the troposphere. Organic peroxy radicals (RO$_2$) are intermediates in the oxidation of hydrocarbons and are precursors for tropospheric ozone formation. This work provides experimental evidence for an increased self-reaction rate of the organic peroxy radical, hydroxyethyl peroxy radical (β-HEP; HOCH$_2$CH$_2$O$_2$), in the presence of water vapor. β-HEP is a derivative of ethylene and plays a significant role in tropospheric chemistry. Ethylene is released both naturally from vegetation and from combustion of wood and fossil fuels with an estimated total emission of 18–25 Tg yr$^{-1}$. A typical urban environment contains 50 µg m$^{-3}$ of ethylene with peak concentrations up to 1000 µg m$^{-3}$ during high traffic hours, whereas rural environments range from 1–5 µg m$^{-3}$. As the troposphere continues to get warmer and wetter, more complexes with water will form, which in turn may perturb the kinetics and product branching ratios of atmospherically important reactions. For example, all unsaturated hydrocarbons will add OH across the double bond, and the addition of O$_2$ will produce β-hydroxy peroxy radicals. Isoprenes and terpenes are two major examples of biogenic sources that will produce β-hydroxy peroxy radicals. If β-HEP demonstrates a self-reaction rate enhancement, we hypothesize that peroxy radicals formed from other unsaturated hydrocarbons will also show enhanced rates with water vapor.
In 1975, Hamilton\textsuperscript{27} discovered that water vapor increases the self-reaction rate of hydroperoxy radical (HO\textsubscript{2}). Since then, numerous studies have verified and refined the findings of this first study.\textsuperscript{28-32} The complete rate expression, including the temperature and pressure dependence of this reaction, was published in 1984.\textsuperscript{30} More recent studies have evaluated the equilibrium constant for formation of the HO\textsubscript{2}-H\textsubscript{2}O complex as well as the rate enhancement by NH\textsubscript{3} and CH\textsubscript{3}OH.\textsuperscript{33} Hamilton’s original work stated that enhancement of the HO\textsubscript{2} self-reaction rate occurs with polar gases and hypothesized that the increased reaction rate was caused by formation of a complex between HO\textsubscript{2} and H\textsubscript{2}O. Existence of an HO\textsubscript{2}-H\textsubscript{2}O complex has since been experimentally verified by Suma \textit{et al.}\textsuperscript{34}

The catalytic effect of water has been discussed at length without agreement on a general theory for predicting a change in the rate of specific reactions in response to changes in water vapor concentration.\textsuperscript{35-37} Stockwell\textsuperscript{38} modeled the water vapor enhancement of the HO\textsubscript{2} self-reaction as it affects atmospheric HO\textsubscript{x} and O\textsubscript{3} budgets, and predicted decreased concentrations of HO\textsubscript{2} and H\textsubscript{2}O\textsubscript{2} under saturated water conditions and thereby decreasing the O\textsubscript{3} concentration. The concentrations of O\textsubscript{3} and HO\textsubscript{x} were previously overestimated by as much as 75% at the surface and 30% at 10 km. Butkovskaya \textit{et al.}\textsuperscript{39} recently showed that 50% relative humidity (4×10\textsuperscript{17} molecules cm\textsuperscript{-3} at 298 K), displaced the branching ratio of the reaction of HO\textsubscript{2} + NO away from NO\textsubscript{2} production and increased the production of HNO\textsubscript{3} by 800%. They hypothesized that a water molecule complexed to HO\textsubscript{2} serves as an energy sink and consequently drives formation of HNO\textsubscript{3} at the expense of OH and NO\textsubscript{2}.

Although the HO\textsubscript{2} self-reaction rate is increased by water, water has no measurable effect on the self-reaction rate of methyl peroxo radical (CH\textsubscript{3}O\textsubscript{2}) or the reaction of CH\textsubscript{3}O\textsubscript{2} + HO\textsubscript{2}.\textsuperscript{37} The absence of catalysis by water is explained by the small fraction of CH\textsubscript{3}O\textsubscript{2} complexed with water.
owing to the small binding energy, calculated to be 8.8 kJ mol$^{-1}$. At 100% relative humidity and 298 K, the percentage of CH$_3$O$_2$-H$_2$O complex is only 0.02%, substantially less than the 8–15% calculated for the HO$_2$-H$_2$O complex,$^{40}$ which is calculated to be bound by 28.9 kJ mol$^{-1}$. This is comparable to the binding energy calculated for the $\beta$-HEP-H$_2$O complex (27.4 kJ mol$^{-1}$). Recent computational work$^{41}$ indicates that other alkyl peroxy radicals can also complex with water vapor. Khan et al.$^{42}$ in a recent atmospheric global modeling study using the STOCHEM-CRI model, predicted that as much as 17% of peroxy radicals in the upper troposphere can exist as complexes.

$\beta$-HEP serves as a model system for investigating the possible role of water vapor in perturbing the kinetics and product branching ratios of other alkyl peroxy radicals. For water vapor enhancement to be observed, a significant portion of reactants must form a radical-molecule complex.$^{43}$ $\beta$-HEP is chosen as a model radical to test the theory that strongly hydrogen-bonded complexes increase the rate of self-reaction. Clark et al.$^{41}$ demonstrated that the optimized structure predicted for the $\beta$-HEP-H$_2$O complex is stabilized by the formation of two intermolecular hydrogen bonds, as is the HO$_2$-H$_2$O complex. The self-reaction rate constant of $\beta$-HEP has previously been established in several studies.$^{44-49}$ However, the kinetics of the effects of water vapor on the self-reaction of $\beta$-HEP has not been studied. Khan et al.$^{42}$ concluded in their modeling study “the reaction between RO$_2$.H$_2$O and NO, RO$_2$/HO$_2$ are worthy of investigation.”

This study presents experimental evidence for the catalytic effect of water vapor on the $\beta$-HEP self-reaction. We derive an empirical equation from experimental measurements expressing the $\beta$-HEP self-reaction rate constant as a function of temperature and water vapor. High level ab initio calculations are used to calculate the equilibrium constant for the formation of the $\beta$-HEP-
H$_2$O complex. The temperature range and water vapor concentrations were chosen because of their significance to conditions in the troposphere.

Figure 2.1- Optimized geometries of β-HEP. The bond length is reported in angstroms.

Figure 2.2- Optimized geometries of β-HEP-H$_2$O complex. Bond lengths are reported in angstroms.
2.4 Methods

2.4.1 Computational Methods

Gaussian 03\textsuperscript{50} was used to compute the equilibrium constant for complexation of β-HEP with H\textsubscript{2}O. Both harmonic and anharmonic vibrational frequencies and rotational constants of both the β-HEP molecule and the β-HEP-H\textsubscript{2}O complex were optimized with the B3LYP/aug-cc-pVDZ method and basis set. The energies were refined by performing single-point energy calculations with the MP2 and CCSD(T) methods. One global minimum and one local minimum structure were identified for the β-HEP molecule. These structures are shown in Figure 2.1. One global and two local minima energy structures for β-HEP-H\textsubscript{2}O (Figure 2.2) complex were identified by Gaussian optimizations through Random Constrained Sampling (RCS) methodology\textsuperscript{36} (see Figure 2.2). The energies of each local and global minimum energy structure were zero-point energy corrected and basis set superposition error corrected\textsuperscript{51} from energies calculated with the B3LYP/aug-cc-pVDZ method and basis set in Gaussian 03. All local minima with energies within $kT$ (2.5 kJ/mol) of the global minimum geometry have measurable probability of being populated at room temperature, and therefore all the geometries were used in the equilibrium constant calculation. Additional details about the partition function calculations and equilibrium constant calculations\textsuperscript{52} are reported in Supplemental material section S-1.

2.4.2 Instrumentation

Figure 2.3 shows the principal components and layout of the experimental apparatus. The photolysis cell is 244 cm long including the end boxes (20 cm $\times$ 18 cm) used to house long-pass Herriott cell optics. The end boxes support the glass photolysis cell and are purged with nitrogen to protect optical surfaces from corrosive gases. The Pyrex cylindrical reaction tube (187 cm in
length, 5.1 cm in diameter) between the end boxes constitutes the effective path length for photolysis and absorption experiments. The reaction tube is surrounded by a Pyrex jacket through which methanol, or a glycol/water mixture is circulated for temperature regulation. All gases are precooled or heated before entering the reaction cell by a jacketed long-path manifold that is either cooled with methanol or heated with a glycol/water mixture. K-type thermocouples placed at both ends of the cell monitor the temperature of the system. Typically, less than 1°C difference occurs between the gases entering and exiting the cell. CaF$_2$ windows are used at both ends of the cell to transmit both UV and IR light over the range of interest. Kinematically mounted mirrors in the optical path of the system make it possible to switch between UV and IR detection methods without changing the position and alignment of the flow cell, light sources and detectors. A purge box placed over the IR and UV optics in the path where

Figure 2.3- Schematic of UV/flash-photolysis system for kinetic measurements.
the excimer-light and D2-light intersect is purged with a steady flow of nitrogen to reduce production of ozone formed by photolysis of O2 outside the cell by the excimer laser. Gas mixtures of N2/O2/2-chloroethanol/H2O flow continuously through calibrated Teledyne-Hastings (HFC series) flow controllers at a total flow rate of 14,000 standard cm³ min⁻¹ (sccm). 2-chloroethanol (99% pure) was purchased from Aldrich and used without further purification. Compressed N2 and O2 cylinders (99%) were obtained from Airgas. Pressure was maintained at 200 Torr total pressure and measured with an MKS-Baratron 690A pressure transducer.

2.5 Measurements

2.5.1 Radical Generation and Measurement

Flash photolysis/UV absorption spectroscopy was used to create and monitor β-HEP. β-HEP was formed in the reaction cell by introducing a gas mixture containing typically $6.1 \times 10^{18}$ molecules cm⁻³ of N2, $6.6 \times 10^{17}$ molecules cm⁻³ of O₂, and $1-6 \times 10^{15}$ molecules cm⁻³ of 2-chloroethanol (HOCH₂CH₂Cl). HOCH₂CH₂Cl was introduced into the gas mixture via dual 100-mL bubblers in series with N₂ at 2000 sccm as the carrier gas. The amount of HOCH₂CH₂Cl is approximate, other gases were metered. Two bubblers in series provide time for carrier N₂ to be saturated with HOCH₂CH₂Cl. The temperature of the bubblers was maintained by immersion in a constant temperature bath maintained at 303 K through use of resistive heaters and a PID controller. Total flow rate was maintained at 14,000 sccm, resulting in a residence time of ~4 s in the reaction cell. A 15 ns pulse from the ArF excimer laser (LasmaPhysik, model EMG201MSC) at 193 nm and energy of 200 to 400 mJ pulse⁻¹ photolyzes an area 2.5 cm wide by 1 cm high. The excimer was fired at 0.33 Hz, allowing for a mostly fresh (~95%) mixture of gas to be probed with every laser shot. The pulse photolyzed a small fraction of the HOCH₂CH₂Cl
molecules to produce HOCH$_2$CH$_2$ and Cl radicals.$^{48, 53}$ Subsequent reaction of HOCH$_2$CH$_2$ radical with O$_2$ produced the β-HEP radical. The initial measured concentration of β-HEP radical is 3–6×10$^{13}$ molecules cm$^{-3}$ depending on the concentration of HOCH$_2$CH$_2$Cl and the intensity of the excimer laser pulse. Along with β-HEP formed by the cleavage of the C-Cl bond in HOCH$_2$CH$_2$Cl, trace amounts of derivatives of β-HEP are made by dissociation of a C-H or C-O bond in HOCH$_2$CH$_2$Cl. Based on the bond dissociation energies for C-H, C-O and C-Cl bonds in HOCH$_2$CH$_2$Cl, other peroxy radicals including ClCH(O$_2$)CH$_2$OH, ClCH$_2$CH(O$_2$)COH and ClCH$_2$CH$_2$O$_2$ are < 0.01% of the β-HEP radical concentration.

Time-resolved detection of the β-HEP radical is done by directing the output from a 30-Watt D$_2$ lamp (Oriel model 60010) through the center of the reaction cell. The excimer laser photolysis beam is aligned coaxially with the UV probe beam using dielectric mirrors that reflect a 30-nm wide band centered on the 193 nm photolysis beam while passing all other UV wavelengths. UV light absorption is detected by a monochromator (Andor model SR-303i-B with a 1800 mm$^{-1}$ or 600 mm$^{-1}$ grating, depending on the light intensity, which decreases with decreasing wavelength in the D$_2$ lamp), and either an intensified and gated CCD camera (Andor model CCD30-11) or a photomultiplier tube (PMT Sens-Tech P30A-14).

The CCD was used in exploratory experiments to identify the major species present. The CCD measured absorption over a wide wavelength range, from 220 to 350 nm, enabling identification of secondary chemistry products. Absorption data at many wavelengths permits a least-squares fit to determine multiple species present in the reaction cell. Time resolved data were collected with the CCD by varying the delay from 1 to 10 ms after the excimer laser was fired. The CCD integrated a 5 µs window after the laser was fired, and 500 shots of the laser were averaged to achieve a sufficient signal-to-noise ratio. These experiments are time-
consuming because measurement of a decay of β-HEP requires a minimum of 30 delay times to achieve a reproducible fit. At 0.33 Hz, 12 h are needed to obtain a single fit to the self-reaction decay.

Because of the fast response, the kinetic data for the β-HEP self-reaction reported here were collected exclusively with the PMT. The PMT together with a monochromator and notch filters with selected wavelengths of 220, 230, and 254 nm (Figure 2.4 adapted from previously published data\textsuperscript{47, 54, 55}) were used to collect kinetic data. The PMT collected data at 30 MHz after the excimer laser was fired. Deconvolution of species concentrations from the total absorption data collected by the PMT requires a decay curve to be collected at each of the three wavelengths. Sufficient signal to noise was obtained by averaging 300–400 shots of the excimer laser. Dark counts were recorded before each shot, and \( I_0 \) was taken as the signal prior to the next laser pulse. Collection of dark counts, blank \((I_0)\) and response \((I)\) was controlled through a LabVIEW routine written in-house.

After collecting decay curves at each of the three wavelengths, recommended cross sections were used from the MPI-Mainz UV-VIS spectral atlas of gaseous molecules\textsuperscript{54} to determine concentrations of \( O_3 \), \( HO_2 \) and β-HEP by a least-squares fit (See supplemental material Section S-2). The β-HEP self-reaction rate constants and β-HEP concentration at time zero were obtained by fitting time-dependent concentrations to a system of differential equations defined by the reactions listed in Table 1.

2.5.2 Water Vapor Generation and Measurement

Water vapor was introduced into the reaction cell by passing \( N_2 \) carrier gas through a bubbler immersed in a constant temperature bath. The amount of water vapor in the reaction cell was controlled by both the temperature of the water and the percentage of total flow of the
carrier gas passing through the bubbler. The water vapor concentration was varied between 
$1.0 \times 10^{15}$ and $2.3 \times 10^{17}$ molecules cm$^{-3}$. Water vapor concentration was held constant during each experiment, varying no more than $\pm 8\%$ over the 8 h duration of an experiment.

Wavelength modulated IR diode laser spectroscopy was used to quantify the water vapor in the cell. Figure 2.3 shows the IR system components. IR light was scanned over a narrow frequency band (100 Hz) centered at the 6790.65 cm$^{-1}$ line such that the entire peak of the selected rovibrational transition of water vapor was observed. The IR beam was produced by a New Focus TLB-6326 tunable diode laser with a line width of $< 300$ kHz. With the use of Herriott cell optics, the IR beam was passed five times through the reaction cell and finally directed towards an IR 1-GHz low-noise photoreceiver (New Focus 1611). The signal from the detector was synchronized with the 100 Hz modulation frequency and visualized and recorded by a digital oscilloscope (TeKtronix model TDS 3024B). Data were processed and converted from absorption measurements to concentrations with a LabVIEW routine. Absorption line strengths were obtained from the HITRAN$^{56}$ database and converted to cross sections.

Figure 2.4- Plot of absorption cross sections of $\beta$-HEP, O3, and HO2 used to fit absorption data to obtain the concentration of each species. The bars superimposed over the plot are bandpass filters used to select the specific wavelengths used for data analysis.
2.5.3 β-HEP Self-Reaction Kinetics Measurements

The experimental set-up was tested and calibrated by measuring the β–HEP self-reaction rate in the absence of water vapor. The dry rate constant measured was $(2.25 \pm 0.15) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ at 296 K, which is within the experimental uncertainty of the recommended NIST evaluation value, $(2.31 \pm 0.3) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ (see Figure 2.5).

The self-reaction of β-HEP follows two product branches, reactions 1 and 2

\[ \text{HOCH}_2\text{CH}_2\text{O}_2 + \text{HOCH}_2\text{CH}_2\text{O}_2 \rightarrow \text{2 HOCH}_2\text{CH}_2\text{O} + \text{O}_2 \quad (2.1) \]

\[ \text{HOCH}_2\text{CH}_2\text{O}_2 + \text{HOCH}_2\text{CH}_2\text{O}_2 \rightarrow \text{HOCH}_2\text{CH}_2\text{OH} + \text{HOCH}_2\text{CHO} + \text{O}_2 \quad (2.2) \]

The overall loss of β-HEP from both reactions was measured in these experiments. The observed rate constant, \( k_{\text{obs}} \), is the overall rate constant representing the following four elementary reactions:

\[ \text{HOCH}_2\text{CH}_2\text{O}_2 + \text{HOCH}_2\text{CH}_2\text{O}_2 \rightarrow \text{Products} \quad (2.3) \]

\[ \text{HOCH}_2\text{CH}_2\text{O}_2 + \text{H}_2\text{O} = \text{HOCH}_2\text{CH}_2\text{O}_2\text{-H}_2\text{O} \quad (2.4) \]

\[ \text{HOCH}_2\text{CH}_2\text{O}_2\text{-H}_2\text{O} + \text{HOCH}_2\text{CH}_2\text{O}_2 \rightarrow \text{Products} \quad (2.5) \]

\[ \text{HOCH}_2\text{CH}_2\text{O}_2\text{-H}_2\text{O} + \text{HOCH}_2\text{CH}_2\text{O}_2\text{-H}_2\text{O} \rightarrow \text{Products} \quad (2.6) \]
193 nm light was used to initiate production of β-HEP, but O₂ is also photolyzed to form singlet oxygen atoms that combine with O₂ to form ozone. The absorption cross section of ozone overlaps strongly with the cross section of β-HEP and HO₂ (Figure 2.4)(postulated to be formed due to undesirable secondary chemistry), causing an interference with direct absorption measurements. However, O₂ could not be excluded from the reaction gas mixture because O₂ is essential for the production of β-HEP and is used to convert HOCH₂CH₂ to β-HEP. Smaller flow rates of O₂ into the reaction mixture lengthened the conversion time of HOCH₂CH₂ into β-HEP, and unwanted secondary chemistry of HOCH₂CH₂ increased. Higher concentrations of O₂ in the reaction cell led to increased production of O₃ that complicated the analysis due to the overlapping absorption spectrum of O₃ with β-HEP. To reduce ozone formation and minimize

Table 1- Reactions used to fit β-HEP decays to obtain the self-reaction rate constant

<table>
<thead>
<tr>
<th>Number</th>
<th>Reaction</th>
<th>$k$ (cm² molecule⁻¹ s⁻¹)</th>
<th>Percent difference in $k_{obs}$ and $k$</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.1</td>
<td>HOCH₂CH₂O₂ + HOCH₂CH₂O₂→ 2 HOCH₂CH₂O + O₂</td>
<td>from fit</td>
<td>0</td>
<td>–</td>
</tr>
<tr>
<td>1.2</td>
<td>HOCH₂CH₂O₂ + HOCH₂CH₂O₂→ HOCH₂CH₂OH + HOCH₂CHO + O₂</td>
<td>from fit</td>
<td>0</td>
<td>–</td>
</tr>
<tr>
<td>1.3</td>
<td>HOCH₂CH₂O₂ + HO₂→ HOCH₂CH₂OO + O₂</td>
<td>$1.0 \times 10^{-11}$</td>
<td>0</td>
<td>[36]</td>
</tr>
<tr>
<td>1.4</td>
<td>HOCH₂CH₂O + O₂→ HOCH₂CHO + HO₂</td>
<td>$5.0 \times 10^{-26} e^{-(-34.92kJ/R)}$</td>
<td>0</td>
<td>[38]</td>
</tr>
<tr>
<td>1.5</td>
<td>HO₃ + HO₂→ H₂O₃ + O₂</td>
<td>$2.2 \times 10^{-13} e^{-(-5.99kJ/R)}$</td>
<td>0</td>
<td>[38]</td>
</tr>
<tr>
<td>1.6</td>
<td>HO₃ + O₂→ CH₂O + HO₂</td>
<td>$9.6 \times 10^{-12}$</td>
<td>0</td>
<td>[25]</td>
</tr>
<tr>
<td>1.7</td>
<td>HO₂ + O₃→ HO + 2O₂</td>
<td>$1.1 \times 10^{-14} e^{-(-4.16kJ/R)}$</td>
<td>0</td>
<td>[13]</td>
</tr>
<tr>
<td>1.8</td>
<td>HO₂ + Cl→ O₂ + HCl</td>
<td>$1.8 \times 10^{-14} e^{-(-4.41kJ/R)}$</td>
<td>0</td>
<td>[13]</td>
</tr>
<tr>
<td>1.9</td>
<td>Cl + O₃→ ClO + O₂</td>
<td>$2.8 \times 10^{-11} e^{-(-2.30kJ/R)}$</td>
<td>±1.1</td>
<td>[13]</td>
</tr>
<tr>
<td>1.10</td>
<td>O + O₂→ O₃</td>
<td>$2.81 \times 10^{-12}$</td>
<td>±0.3</td>
<td>[13]</td>
</tr>
<tr>
<td>1.11</td>
<td>Cl + CICH₂CH₂OH→ OCH₂CH₂O₂ + HCl</td>
<td>$3.08 \times 10^{-11}$</td>
<td>±0.9</td>
<td>[38]*</td>
</tr>
<tr>
<td>1.12</td>
<td>Cl + CICH₂CH₂OH→ HO₂ + HCl + CH₂CHO</td>
<td>$3.08 \times 10^{-11}$</td>
<td>±0.2</td>
<td>[38]*</td>
</tr>
<tr>
<td>1.13</td>
<td>HO₂ + OH→ H₂O + O₂</td>
<td>$4.8 \times 10^{-11} e^{-(-0.88kJ/R)}$</td>
<td>0</td>
<td>[13]</td>
</tr>
<tr>
<td>1.14</td>
<td>OH + O→ O₂ + H</td>
<td>$3.59 \times 10^{-11}$</td>
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<td>[13]</td>
</tr>
<tr>
<td>1.15</td>
<td>Cl + Cl→ Cl₂</td>
<td>$1.46 \times 10^{-13}$</td>
<td>0</td>
<td>[13]</td>
</tr>
<tr>
<td>1.16</td>
<td>O + HO₂→ OH + O₂</td>
<td>$5.8 \times 10^{-11}$</td>
<td>0</td>
<td>[13]</td>
</tr>
<tr>
<td>1.17</td>
<td>OH + ClO→ ClOOH</td>
<td>$1.46 \times 10^{-11}$</td>
<td>0</td>
<td>[38]</td>
</tr>
<tr>
<td>1.18</td>
<td>H + O₂ + M→ HO₂ + M</td>
<td>$5.4 \times 10^{-22}$</td>
<td>0</td>
<td>[38]</td>
</tr>
<tr>
<td>1.19</td>
<td>H + O₂→ OH + O₂</td>
<td>$2.84 \times 10^{-11}$</td>
<td>0</td>
<td>[38]</td>
</tr>
<tr>
<td>1.20</td>
<td>Cl + O₂→ ClO</td>
<td>$9.11 \times 10^{-15}$</td>
<td>0</td>
<td>[13]</td>
</tr>
</tbody>
</table>

*Analogous rate constant of CH₃CH₂OH used for CICH₂CH₂OH.
unknown chemistry between $O(^1D)$ radicals and other species while still producing $\beta$-HEP in less than 30 $\mu$s, the $O_2$ concentration was decreased until the $\beta$-HEP concentration was 3-6 times the detection limit. The ozone concentration was between $3-5 \times 10^{16}$ molecules cm$^{-3}$ in the experiment.

Detection of $\beta$-HEP decays was begun 1 ms after the photolysis laser fired because residual reflections from the photolysis pulse interfered with detection by the PMT. Because photolysis of HOCH$_2$CH$_2$Cl can form HO$_2$, absorption of HO$_2$ was also considered. Three collection wavelengths at 220, 230, and 254 nm provide sufficient data for a least-squares fit to obtain the concentrations of O$_3$, HO$_2$, and $\beta$-HEP (see supplemental material section S2 for the equations used for the least squares fit and Figure 2.6A and 2.6B for typical decay curves). At each condition, multiple decay curves were measured until $k_{obs}$ was within the 95% confidence interval of the mean. Typically, 9 to 34 $\beta$-HEP decay curves were collected at different initial $\beta$-HEP concentrations at each water vapor and temperature condition. The HO$_2$ concentration was always found to be below the detection limit ($\sim 1 \times 10^{13}$ molecules cm$^{-3}$)(i.e. the detection limit of the CCD and PMT was 0.006 absorbance units, therefore the detection limit for [HO$_2$] at 220nm=[0.006 $\div$ (3.16$\times$10$^{-18}$ cm$^2$/molecule $\times$ 187 cm)]=1$\times$10$^{13}$ molecule/cm$^3$), and thus the least-squares fit was reduced to a two-parameter least-squares fit (Equation 1) Derivation of Equation 2-1 is shown in supplemental material section S-2

$$A_x = b \left( \sigma_x^{O_3}[O_3] + \sigma_x^{\beta-HEP}[\beta \text{ - HEP}] \right)$$  \hspace{1cm} \text{Equation (2.1)}
Where b=pathlength, \( \sigma \)=cross section, A=absorbance and \( x=220, 234 \) and 254nm.

Figure 2.5- Comparison of the measured \( \beta \)-HEP self-reaction rate constant in the absence of water vapor as a function of temperature with the Atkinson et al. (NIST) published values. NIST values represented as solid line in the figure.

Figure 2.6- (A) Representative data collected at the three wavelengths 220 (green), 230 (red), and 254 (black). (B) A typical fit to the measured decay.
2.6 Results

2.6.1 β-HOCH₂CH₂O₂ Self-Reaction Kinetics

Each curve was fit to a model composed of coupled differential equations describing the reactions given in Table 1. Measured time-dependent concentrations of β-HEP, and calculated concentrations (using the photon flux and quantum yields of O₂ and H₂O photolysis) of oxygen atoms, hydrogen atoms and hydroxy radicals are inputs to the model. The outputs are concentration of β-HEP at time 0, the observed self-reaction rate of β-HEP (kₜₒₙₜ), and the HO₂ concentration. Table 2 presents the results of all experiments with 95% confidence intervals for [H₂O], [β-HEP]₀, and kₜₒₜ at 274 K, 280 K, 288 K, and 296 K. Figure 2.6a shows representative decays collected at 220, 230 and 254 nm. Figure 2.6b is a typical fit to the measured β-HEP decay. The total experimental error in kₜₒₜ ranges from 4 to 10% depending on experimental temperature and concentrations of O₃ and water vapor. This error could be reduced by decreasing the amount

<table>
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<tr>
<th>Temperature (K)</th>
<th>[H₂O] (molecules cm⁻³ / 1 x 10¹⁶)</th>
<th>kₜₒₜ (fit) (cm³ molecule⁻¹ s⁻¹)</th>
<th>[β-HEP]₀ (molecules cm⁻³ / 1 x 10¹¹)</th>
<th>kₜ (measured) (cm² molecule⁻¹ s⁻¹)</th>
<th>Number of Trials</th>
</tr>
</thead>
<tbody>
<tr>
<td>274</td>
<td>0.18 ± 0.08</td>
<td>3.77 ± 0.28</td>
<td>3.4 ± 0.1</td>
<td>2.82 ± 0.03</td>
<td>21</td>
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<tr>
<td></td>
<td>1.57 ± 0.14</td>
<td>6.04 ± 0.48</td>
<td>4.5 ± 0.2</td>
<td></td>
<td>12</td>
</tr>
<tr>
<td></td>
<td>2.60 ± 0.13</td>
<td>8.14 ± 0.68</td>
<td>3.1 ± 0.2</td>
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<tr>
<td></td>
<td>3.32 ± 0.33</td>
<td>8.54 ± 0.71</td>
<td>2.9 ± 0.1</td>
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<tr>
<td>280</td>
<td>0.29 ± 0.13</td>
<td>4.16 ± 0.50</td>
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<td>3.00 ± 0.13</td>
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<td></td>
<td>2.10 ± 0.08</td>
<td>5.36 ± 0.42</td>
<td>4.0 ± 0.3</td>
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<tr>
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<td>4.70 ± 0.25</td>
<td>5.70 ± 0.50</td>
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<td>5.86 ± 0.16</td>
<td>6.68 ± 0.67</td>
<td>3.6 ± 0.6</td>
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<td>6.76 ± 0.23</td>
<td>8.47 ± 0.97</td>
<td>3.1 ± 0.3</td>
<td></td>
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<tr>
<td>288</td>
<td>0.20 ± 0.05</td>
<td>3.38 ± 0.17</td>
<td>3.9 ± 0.1</td>
<td>2.5 ± 0.02</td>
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<td>3.80 ± 0.10</td>
<td>4.09 ± 0.28</td>
<td>4.1 ± 0.7</td>
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<td>8.21 ± 0.29</td>
<td>4.94 ± 0.28</td>
<td>4.7 ± 0.9</td>
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<td>12.20 ± 0.07</td>
<td>6.20 ± 0.57</td>
<td>5.0 ± 0.6</td>
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<tr>
<td>295</td>
<td>0.18 ± 0.05</td>
<td>2.92 ± 0.18</td>
<td>3.1 ± 0.1</td>
<td>2.25 ± 0.15</td>
<td>34</td>
</tr>
<tr>
<td></td>
<td>6.26 ± 0.20</td>
<td>3.67 ± 0.45</td>
<td>5.1 ± 0.3</td>
<td></td>
<td>16</td>
</tr>
<tr>
<td></td>
<td>11.80 ± 0.03</td>
<td>3.54 ± 0.46</td>
<td>5.3 ± 0.2</td>
<td></td>
<td>10</td>
</tr>
<tr>
<td></td>
<td>22.50 ± 0.07</td>
<td>4.16 ± 0.14</td>
<td>5.9 ± 0.3</td>
<td></td>
<td>12</td>
</tr>
</tbody>
</table>
of O$_2$ in the cell, but, as noted previously, reducing oxygen would not produce sufficient β-HEP.

Sensitivity of the kinetic model was done to determine the effects of uncertainties in the rate constants on fitting $k_{obs}$. The rate constants were changed individually by ±30%, and the measured β-HEP decay curve were fit to obtain an adjusted value for $k_{obs}$, i.e. $k'$. This $k'$ was compared against the $k_{obs}$ obtained from the fit of the experimental data with the original model. These results are listed in Table 1 as the percent difference between $k_{obs}$ and $k'$. For example, a 30% change in the rate constant of the O$_3$ + Cl reaction leads to a 1% change in calculated $k_{obs}$. The uncertainty in $k_{obs}$ caused by uncertainties in the reactions of the model is estimated to be 2.5% (see table 1). The sensitivity analysis shows that secondary chemistry is dominated by reactions involving O$_3$ and Cl radicals.

![Graph](image)

Figure 2.7- Observed rate constant ($k_{obs}$) for self-reaction of β-HEP versus water vapor concentration. The lines represents the fit to the rate expression (Eq. (III)) describing the rate as a function of both water vapor concentration and temperature.
In their HO₂ and HO₂-H₂O study, Cox and Burrows,²⁸ Kircher and Sander,³⁰ and Lii et al.³¹ applied the assumption that HO₂ and complex absorption cross sections are not resolveable in the UV. But Aliosio et al.,⁵⁸ showed that the complex cross section in the UV is red shifted by 59nm, which contradicts above assumption. Therefore, in this study, we account for both complex and β-HEP cross section by assuming the following relationship,

\[
\sigma_x^{β-HEP-H_2O} = f_x \sigma_x^{β-HEP}
\]

where \( x = 220, 230 \) and \( 254 \) and \( f_x \geq 0 \).

An expression that describes the water vapor dependence of the β-HEP self-reaction must include the reactions of β-HEP with both β-HEP and β-HEP-H₂O and the reaction of β-HEP-H₂O with β-HEP-H₂O. A reduced form of the rate law is

\[
k_{obs} = k_3 + k_5[H_2O]
\]

Equation (2.2)

where \( k_3 \) and \( k_5 \) are rate constants for reactions 3 and 5 respectively, and \( K \) is the equilibrium constant for reaction 4 (See supplemental material Section S-3 for a derivation of Equation 2). The same expression has previously been used to describe water vapor enhancement of the HO₂ self reaction. Equation 3, which is a functional form of equation 2 was used to obtain the temperature dependence of \( k_{obs} \) (see supplemental material section S-4 for the derivation)

\[
k_{obs} = k_3 \left[1 + A[H_2O] \times \exp(-E/RT)\right]
\]

Equation (2.3)
The measured $k_{obs}$ were fit with Equation 2.3, which includes the contribution from the dry rate constant ($k_3$), the water vapor dependence, and the temperature dependence of the product, $k_5K$.

Fits to the data in Figure 2.7 with Equation 2.3 are shown as solid lines. The enhancement from water vapor in Equation 3 is taken from the data collected in this work. The dry rate constant ($k_3$) values from this work are at 4 different temperatures only, therefore we have used the dry rate constant equation from Atkinson's (NIST) evaluation of all previously available data on dry rates. Our data agree with the NIST evaluation (see Figure 2.5). Kircher included a pressure dependent term in the model for the HO$_2$ self-reaction but no data for the pressure dependence of the β-HEP self-reaction is currently available. The best fit to the data in figure 2.7 yields the following rate expression:

Figure 2.8- Effect of temperature on K and wet rate constant $k_5$. The solid line represents the change in the wet rate constant $k_5$, and the dotted line (---*) represents the change in K. The dashed lines represent the uncertainty in K and $k_5$. The highlighted area represents the range of $k_5$ values limited by the hard sphere collision rate constant. Atkinson (NIST) evaluation for the dry rate constant (---**--) as a function temperature.
In their investigation of the kinetics of the HO$_2$ + HO$_2$-H$_2$O reaction, Kanno et al. measured $K$ and the rate constants independently. In this study, the product of $K$ and $k_5$ is defined by the slopes of plots of $k_{obs}$ vs [H$_2$O] at each temperature (see Equation 2.2 and Figure 2.7), and $K$ and $k_5$ are inseparable from the data collected. Therefore, the equilibrium constant for the formation of the $\beta$-HEP-H$_2$O complex ($K$) was estimated from $ab$ initio calculations, as described previously in the computational method section 2.1. The calculated $K$ values as a function of temperature are given in Figure 2.8. The dashed lines around $K$ in figure 2.8 are the uncertainty limits in $K$ assuming an uncertainty in $\beta$-HEP-H$_2$O complex binding energy of $\pm 6.3$ kJ mol$^{-1}$.

Dividing the slopes ($k_5K$) of the fits in figure 2.7 by the calculated $K$ yields $k_5$ values as shown in Figure 2.8. The uncertainty in $K$ propagates into the calculation of $k_5$, shown as dashed lines in Figure 2.8. Since the calculated $k_5$ values cannot be larger than the hard sphere collision rate constant shown by the solid line in Figure 2.8, the calculated $K$ values must be too small. Thus, the true $K$ values must be near the upper uncertainty limit, and consequently, the $k_5$ values must be near the lower uncertainty limit in Figure 2.8. Kanno et al. also found the $ab$ initio calculated $K$ values for the HO$_2$-H$_2$O complex to be an order of magnitude smaller than their experimentally determined $K$ values. This suggests that $ab$ initio calculated values of $K$ may in general be an order of magnitude smaller than measured values of $K$ for formation of radical-water complexes. In any case, the rate constant for $\beta$-HEP + $\beta$-HEP-H$_2$O ($k_5$) is 1 to 3 orders of

$$k_{obs}(T,[H_2O]) = 7.8 \times 10^{-14} \exp((8.3 \pm 2.5\text{kJ/mol})/RT) \times \left(1 + (1.7 \pm 0.2) \times 10^{-30} \times [H_2O] \times \exp((71 \pm 17\text{kJ/mol})/RT)\right)$$

$$= 7.8 \times 10^{-14} \times \exp((8.3 \pm 2.5\text{kJ/mol})/RT) + (13.2 \pm 1.56) \times 10^{-44} \times \exp((79.3 \pm 17.18\text{kJ/mol})/RT) \times [H_2O]$$

Equation (2.4)
magnitude larger than that for $\beta$-HEP + $\beta$-HEP (See Figure 2.8). To better estimate the effects of water vapor on the $\beta$-HEP self-reaction, the equilibrium constant $K$ needs to be determined experimentally, but this is beyond the scope of this paper.

### 2.6.2 HO$_2$ Production and Secondary Chemistry

The formation of some HO$_2$ and Cl radicals in our experiments is inevitable. Determining the concentration of HO$_2$ and the secondary chemistry occurring due to HO$_2$ and Cl radicals is thus necessary. HOCH$_2$CH$_2$Cl has 3 non-equivalent hydrogen atoms, abstraction of either alcohol hydrogen or a hydrogen alpha to the OH group leads to production of HO$_2$.

Extraction of the alcohol hydrogen from HOCH$_2$CH$_2$Cl followed by reaction with O$_2$ leads to HO$_2$ formation via the following reactions:

$$\text{HOCH}_2\text{CH}_2\text{Cl} + \text{Cl} \rightarrow \text{OCH}_2\text{CH}_2\text{Cl} + \text{HCl} \quad (2.7)$$

$$\text{OCH}_2\text{CH}_2\text{Cl} + \text{O}_2 \rightarrow \text{HO}_2 + \text{ClCH}_2\text{CHO} \quad (2.8)$$

The fate of the OCH$_2$CH$_2$Cl radical has not been investigated in the presence of oxygen, but reactivity of the analogous OCH$_2$CH$_3$ radical in air has been reported by Atkinson et al.$^{49}$ They reported a rate constant of $1.1 \times 10^{-14}$ cm$^3$ molecule$^{-1}$ s$^{-1}$ for extraction of the alpha hydrogen by molecular oxygen.

Extraction of an alpha hydrogen by a Cl atom followed by reaction with molecular oxygen also produces HO$_2$ radical via the following reactions:
Attack of a chlorine atom at a beta hydrogen produces 2-chloro-hydroxyethyl peroxy radical.

\[
\text{HOCH}_2\text{CH}_2\text{Cl} + \text{Cl} \rightarrow \text{HOCHCH}_2\text{Cl} + \text{HCl} \quad (2.9)
\]

\[
\text{HOCHCH}_2\text{Cl} + \text{O}_2 \rightarrow \text{HO}_2 + \text{ClCH}_2\text{CHO} \quad (2.10)
\]

To address the question of the relative significance of these reactions, the product branching ratio for extraction of the alcohol, alpha and beta hydrogen was calculated using high level \textit{ab initio} calculations. A slice of the potential energy surface is included in supplementary material Figure S1. Extraction of an alpha hydrogen from HOCH\textsubscript{2}CH\textsubscript{2}Cl has the lowest energy barrier (19 \text{kJ mol}^{-1}) followed in energy by extraction of the beta hydrogen (37 \text{kJ mol}^{-1}) and finally the alcohol hydrogen (100 \text{kJ mol}^{-1}). The product branching ratio and the rate constant reported by Garzon \textit{et al}\textsuperscript{60} and Taatjes \textit{et al}\textsuperscript{61} for Cl + HOCH\textsubscript{2}CH\textsubscript{3} were used as a surrogate for Cl + HOCH\textsubscript{2}CH\textsubscript{2}Cl owing to the similar activation barriers reported (7, 18, 111 \text{kJ mol}^{-1}). The kinetic model in Table 1 includes the product branching ratios for the reaction of Cl + HOCH\textsubscript{2}CH\textsubscript{2}Cl and the other possible secondary chemistry reactions of Cl and HO\textsubscript{2} radical that are listed in Table 1. The model predicts the amount of HO\textsubscript{2} formed in the reaction mixture and the amount of Cl atoms that are lost as a function of time.
Reactions 1.15 and 1.8 in Table 1 both have large rate constants, but make a negligible contribution to the loss of chlorine atoms because the concentrations of Cl and HO₂ are very small. Reaction 1.20 is a three-body reaction, but because of the relative high concentration of O₂ and the total pressure in the cell, accounts for 3% of the loss of chlorine atoms. Aside from the reactions of Cl with HOCH₂CH₂Cl (reactions 1.11 and 1.12) which account for 52% of Cl atom loss, reaction 1.9 accounts for the largest loss of Cl atoms (44%).

The effect of H and OH radicals on the HO₂ concentration are also included in the model. In this study, the 193 nm light was used to initiate photolysis of HOCH₂CH₂Cl in the cell. This wavelength of light is capable of photolyzing both H₂O and O₂ in the cell, leading to formation of H, OH and O radicals. The photon flux and quantum yields for photolysis of H₂O and O₂ predict the formation of 4.5×10¹³ molecules cm⁻³ of H and OH radicals from H₂O photolysis and 3.6×10¹⁶ molecules cm⁻³ of O atoms from photolysis of molecular oxygen. These values were included as initial concentrations of H, OH and O in the model. HO₂ is formed in the cell from reactions 1.4, 1.6, 1.12, and 1.18. H and OH radicals have opposing effects on the HO₂ radical concentration. H radical increases the concentration of HO₂ due to the reaction of H and O₂, whereas OH radical remove HO₂ due to the fast reaction between OH and HO₂. H radical is responsible for formation of 34% of HO₂ radical in the cell, whereas OH radical is responsible for 32% of removal of HO₂ radical. Overall, these effects increase formation of HO₂ radical. The model predicts HO₂ concentration to steadily increase after photolysis. But, at 10 ms, the HO₂ concentration is predicted to be 6×10¹² molecules cm⁻³ and below our detection limit of 1×10¹³ molecules cm⁻³ confirming the experimental observations. HO₂ is not produced in measurable concentration, but can still contribute to the decay of β-HEP and for this reason is included in the model fits.
2.7 Discussion

The current study showed no evidence for significant HO$_2$ production, but in their investigation of β-HEP self-reaction kinetics with a similar source chemistry, Murrells et al.\textsuperscript{48} reported measuring significant concentrations of HO$_2$. They reported initial concentrations of HO$_2$ that were twice that of β-HEP and explained that this was due to an unknown source of HO$_2$ production. One possible explanation for the discrepancy between Murrells et al. and this work is that Murrells et al. did not include O$_3$ in their data analysis and it is likely that the signal attributed to HO$_2$ was instead O$_3$. No concentrations of precursor gas species were reported by Murrells et al. but the published decays of β-HEP have no discernible time delay between the formation of HOCH$_2$CH$_2$ and β-HEP, suggesting a high concentration of O$_2$ was used to drive β-HEP formation. A high concentration of O$_2$ produces a significant amount of O$_3$ from photolysis at 193 nm. In addition, their wavelength selection of 225 and 260 nm complicates the data analysis. The absorption cross-sections of HO$_2$, β-HEP, and O$_3$ are listed in Table 3. At 225 nm the cross sections of HO$_2$ and O$_3$ are the same and the cross section of β-HEP is only 28% greater than 20% of one another, but at 260 nm, the cross section of O$_3$ is 30 times larger than the cross section of HO$_2$, and β-HEP’s cross section is 11 times larger than the cross section of HO$_2$. Thus, the use of 225 and 260 nm as probe wavelengths only provides enough information to deconvolve HO$_2$ and β-HEP from one another. At 225 nm, O$_3$ absorption overwhelms the signal attributed to HO$_2$. The absence of treatment of O$_3$ in Murrells et al.’s work thus explains the discrepancy between the results of Murrells et al. and this work.

Anastasi et al.\textsuperscript{46} reported a β-HEP self-reaction rate constant of $7.69 \pm 1.2 \times 10^{-12}$ cm$^3$ molecule$^{-1}$ s$^{-1}$ at 300K, 3.4 times larger than the dry rate constant, $2.31 \times 10^{-12}$ cm$^3$ molecule$^{-1}$ s$^{-1}$, others reported.\textsuperscript{45, 47, 48, 62} The difference was attributed to the use of different absorption cross
sections for β-HEP\textsuperscript{48}. However, Anastasi \textit{et al.}\textsuperscript{46} used a water vapor concentration of approximately $1 \times 10^{18}$ molecules cm\textsuperscript{-3}, resulting in formation of the reactive β-HEP-H\textsubscript{2}O complex which they did not account for. The rate constant reported by Anastasi \textit{et al.} agrees with the $k_{abs}$ measured in this work (Equation 4).

The data presented in this work describe the relationship between increasing rates of reaction of peroxy radicals and increasing water vapor. \textit{Ab initio} calculations estimated the equilibrium constant to allow determination of both the percentage of β-HEP complexed with water and the rate constant for β-HEP + β-HEP-H\textsubscript{2}O (Reaction 5). Enhancement of peroxy radical self-reaction by water depends on the fraction of radical complexed with water and on the nature of the peroxy radical. For example, HO\textsubscript{2} and β-HEP have equilibrium constants of formation with water within an order of magnitude, but at 50\% relative humidity and 280 K, the rate constant for β-HEP is 6.1 times larger than the dry rate constant whereas the rate constant for HO\textsubscript{2} is 1.5 times larger than the dry rate constant. The calculated fraction of β-HEP radical complexed with water vapor has a large uncertainty mostly due to the uncertainty in the calculation of $K$. Future work must measure the $K$ values experimentally to better estimate the enhancement in the β-HEP self-reaction rate.
The contribution of water vapor to the rate of loss of $\beta$-HEP and HO$_2$ is determined by the ratio $k_{obs}/k_3$ defined here as the enhancement factor. This enhancement factor allows a measure of the contribution of water vapor to the rate of loss of $\beta$-HEP and HO$_2$ at differing temperatures and relative humidities, Figure 2.9. For any given water vapor concentration the rate of loss of $\beta$-HEP is faster than the loss of HO$_2$, e.g. at 298 K the rate of loss of $\beta$-HEP is twice that of HO$_2$.

![Figure 2.9- Enhancement of the self-reaction of $\beta$-HEP and HO2 at various relative humidities and temperatures](image)

2.8 Conclusion

Water vapor catalysis will become increasingly more significant as global warming increases water vapor concentrations in the troposphere. Before the water vapor effect on the HO$_2$ self-reaction was included in atmospheric models, the models predicted a much higher concentration
of HO$_2$ in the atmosphere. As evidenced by the work of Stockwell, inclusion of the effect of water vapor on the self-reaction rate of HO$_2$ leads to a 75% decrease in the prediction of tropospheric HO$_x$ and O$_3$ concentrations. Thus, it is likely that current models overpredict the concentration of β-HEP and other peroxy radicals and consequently overestimate the concentration of O$_3$. β-hydroxy peroxy radicals are atmospherically important because they are products in the reaction of OH radicals with unsaturated VOC’s from diverse emission sources. Isoprene and terpenes are two major examples of biogenic sources that produce β-hydroxy peroxy radicals and should demonstrate a water vapor enhancement of the kinetics. Thus, future studies on the effect of water vapor on the kinetics of peroxy radical reactions should focus on β-hydroxy peroxy radicals and related compounds that form two hydrogen bonds with water and have binding energies greater than 17 kJ mol$^{-1}$. Most likely inclusion of the effects of water vapor on the kinetics of the β-HEP and HO$_2$ self-reactions in atmospheric models will lead to predictions of smaller tropospheric O$_3$ concentrations. To correctly predict pollution levels, atmospheric models also need to be adjusted to include the effects of polar molecules such as methanol and ammonia on the rates of radical reactions.

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3 DISCHARGE-FLOW MASS SPECTROMETER

3.1 Origins of the Flow Tube Method

The discharge-flow mass spectrometer (DFMS) is a great resource in studying the kinetics of atmospheric reactions. The purpose of this section is to provide a brief history of the flow tube with emphasis on the developments that have contributed most to its application as a kinetic tool.

In the 1970s and 1980s, there was an increased demand for gas phase reaction rate data. At the time, there were serious environmental, economic, and social implications of gas phase kinetic data. Models based on these data predicted the impact of anthropogenic chemicals on stratospheric ozone and other key atmospheric reactions. Committees and organizations were formed to collect, evaluate, and disseminate this atmospheric kinetic information. With high concern for accuracy in the atmospheric reaction rate constants, many researchers began to utilize the DFMS as a key source for measurements. During this era, the DFMS served as the most prolific source of atmospheric kinetic data.\(^1\) Figure 3.1 shows a schematic of a simple flow cell that gave rise to today’s DFMS.

The flow tube kinetic method evolved from early discharge tube studies conducted by Wood and Bonhoeffer in the early 1900’s as they experimented with resonance fluorescence in metal atoms.\(^63\) In 1929, Smallwood performed the first flow tube kinetic study, reporting the rate of recombination of hydrogen atoms. Smallwood adapted the Wood-Bonhoeffer type discharge tube in order to create atomic hydrogen,\(^64\) and the atom concentration in the flow tube was measured with a calorimeter attached to the outside of the flow tube.

An important step in DFMS was the invention of the Wrede-Harteck gauge,\(^65\) a simple device for measuring the partial pressure of atoms in a discharged gas mixture. Harteck studied the reactions of atomic oxygen with various compounds. Atomic oxygen was produced by an AC
discharge, its concentration estimated using the Wrede-Harteck gauge, and the reactivity
determined qualitatively by spectroscopic measurements of its chemiluminescence.

In the late 1950’s and 1960’s, several important advances were made in flow tube
techniques and instrumentation. In 1958, Kaufman \(^{66}\) demonstrated that the intensity of the
chemiluminescence was proportional to the concentrations of the reactants, i.e \(I \propto [A]\). This
discovery provided a sensitive and quantitative detection method. Similarly, Kaufman
demonstrated that extremely fast gas phase reactions could be quantified by gas phase titrations.
The concept of gas phase titrations has now been extended to many different radical species, and
serves as a clean, quantitative method in fast stoichiometric reactions.\(^1\)

The use of chemiluminescence as a detection scheme in a flow tube was universal for
several years. Clyne et al.\(^{67}\) were the first to develop UV detectors to monitor atomic oxygen,
and Setser et al.\(^{68}\) reported using an IR detector to study hydrogen extraction. The greatest
development in instrumentation of a flow tube, however, was the implementation of a
quadrupole mass spectrometer as a detector. Foner and Hudson\(^{69}\) and Phillips and Schiff\(^{70}\) were
the first to combine mass spectrometer detectors with flow tubes for kinetic measurements. Since
then, this combination has been successfully exploited in a wide variety of kinetic application,
and resembles mostly to that of today’s DFMS.
3.2 Schematic and Technique of a Kinetic Flow Tube

A schematic of a simple flow tube is shown in Figure 3.1. The major component is a glass tube surrounded by a controlled jacket to maintain steady constant temperature. The inner flow tube surface is often treated with a wall coating to inhibit the removal of radicals on the tube’s surface.

The carrier gas (M) enters the flow tube at the left. This gas is the major component in the flow tube and defines the physical properties of the gas stream, such as pressure, flow velocity, heat capacity, viscosity, etc. This gas also acts as a heat bath to maintain the reactants at the temperature of the glass wall, controls diffusion, and serves as a third body in termolecular reactions. Helium traditionally serves as the carrier gas for its inertness and high thermal conductivity.

Atomic or radical reactants (c) are most commonly generated by a microwave discharge. These reactants can be further converted to different atomic or radical species by a titration reaction before being introduced into the flow cell. The radical reactant enters the flow tube via

Figure 3.1- A simple schematic of a flow tube. (Adapted from Howard)
an inlet downstream from the carrier gas entrance. A small flow of He or Ar is used to flush the
gas through the radical source.

A pressure port is located at the center of the reaction zone to minimize errors and
corrections due to the pressure gradient in the flow tube. The pressure port needs to be oriented
properly so as to measure only the static pressure in the flow tube.

Flow tube kinetic measurements usually are made by varying the reaction zone length ($Z$),
which consists of the distance between the point the two reactants meet and the point of detection.
The reaction zone length can be changed by either moving the point of mixing or the point of
detection. The easiest and most common method is changing the mixing point by using a
moveable injector tube as seen in Figure 3.1. The exit orifice on the moveable injector tube
consists of a series of small holes around the diameter of the tube. The purpose of this
configuration is to provide a rapid mixing of the added gas with the carrier stream. Added
reactant A is assumed to be thoroughly mixed with the carrier gas.

Sensitivity is important when determining the detection method. For the best results, the
concentration of the radical species, $c$, should be less than $10^{11}$ molecules cm$^{-3}$ when radical
reactions are studied. Typically, the concentration of the added reactant, $A$, is much larger by
several orders of magnitude, thus creating a pseudo-first order reaction in $c$.

### 3.3 Calculating the Kinetics in a Flow Tube

When the flow tube is run under pseudo-first order conditions, i.e. $[A] \gg [c]$, the rate
equation is
\[
\frac{dc}{dt} = -k[c] \tag{3.1}
\]

and only the relative concentrations of c need to be measured.

A basic assumption is made that the radical reactant is mixed homogeneously with the carrier gas and that there are no concentration gradients. A second assumption is made that the carrier gas flow velocity (\(v\)) is the same as the radical velocity (usually 300-2000 cm s\(^{-1}\)). With these assumptions, the reaction time (\(t\)) in the flow tube, \(t = Z/v\), is the time from when the radicals and reactants mix at the point of contact to the point of detection. Thus, reaction time and distance are equivalent in the flow system.

A rate constant measurement is made by measuring the radical concentration with the moveable injection tube at several different positions, while the reactant flow rate is held constant (i.e. = 10cm, 20cm, 30cm, etc.). The data are plotted \(\ln[c]\) vs. \(Z\) and the slope is used to calculate the pseudo-first order reaction rate constant

\[
k = -v \frac{d(\ln[c])}{dz} \tag{3.2}
\]

Limitations to this method, i.e. wall loss, axial and radial diffusion, pressure gradients, etc, can all be accounted for and are discussed by Howard.\(^1\) The pseudo-first order rate constant, \(k\), can be converted into a true second order rate constant, \(k'\), by plotting the \(k\) vs. the concentration of the excess reagent. The slope of this line is the true second order rate constant.
3.4 A Modern Discharge Flow Mass Spectrometer

With the addition of the mass spectrometer to the flow tube, further engineering adjustments had to be made to the simple flow tube design shown in Figure 3.1. While flow tubes are normally operated at pressures approximate to 1 Torr, traditional quadrupole mass spectrometers are not able to properly function above \( >10^{-4} \) Torr. In order to couple a flow tube to a quadrupole mass spectrometer, intermediate stages are required to systematically lower the flow tube pressure to ideal quadrupole pressures. Figure 3.2 depicts these intermediate stages, and resembles a modern DFMS system. During the 1970’s and 1980’s, these intermediate stages were pumped down with diffusion pumps, a very messy and expensive method to achieve high vacuum. A majority of the expense came from the constant need of replenishing liquid N\(_2\) in the diffusion pump cold traps. This required 24 hour monitoring and systems in place to refill the liquid N\(_2\) cold traps when necessary.

![Diagram of modern flow tube system](image)

Figure 3.2 - Modern flow tubes require intermediate pump down stages in order achieve pressures suitable for quadrupole mass spectrometry.
As the demand and funds for atmospheric kinetics diminished, many researchers could not afford to maintain the operation of the DFMS systems. In the decades following the 1980’s, these systems were shut-down, stored, or repurposed. By the early 2000’s, only a dozen or so of DFMS systems were operating globally.

3.5 NASA’s Jet Propulsion Laboratory Discharge Flow Mass Spectrometer

Since its inception in 1958, one of NASA’s programmatic goals has been to be the top leader in atmospheric science. During the 1970’s and 1980’s, NASA funded the research of many atmospheric kineticists. A few of these scientists went on to win a Nobel Prize for their research in ozone layer depletion. Many of these researchers employed the use of DFMS in their studies. With time, these researchers could no longer justify the expense of operating the DFMS.

Now, over 30 years later, the need of a functional DFMS is required. Many of the atmospheric rates measured in the 1970’s and 1980’s are in need of reevaluation. In 2014, Cline\textsuperscript{72} demonstrated that atmospheric reaction rates and kinetics can be influenced by the presence of water vapor. With the increase of global temperature, the need to reevaluate atmospheric gas phase reactions as a function of water vapor concentration is required. With this end in mind, Stan Sander of the Jet Propulsion Laboratory and Jaron Hansen of Brigham Young University decided to resurrect an old discharge flow mass spectrometer, and equip it with the most advanced hardware.
Figure 3.3 shows the schematic of the newly engineered DFMS. The initial parts consisted of an inherited two-stage vacuum chamber, a 30 year old Pyrex jacketed flow cell, and a 15 year old mass spectrometer; none-of-which had come from the same system.

The vacuum chamber for the quadrupole mass spectrometer was redesigned to use turbo pumps instead of diffusion pumps to vacuum out the system. The vacuum chamber is a two-stage differentially pumped vacuum system. The first stage is pumped using a Pfeifer Hi Pace 1500 turbo which was backed by an XDS35i scroll pump. The second stage, which houses the quadrupole mass spectrometer, is pumped by a Pfeifer Hi Pace 700 turbo pump, which is backed by a Pfeifer Hi Pace 80 turbo pump, which is then again backed by an XDS15i scroll pump. An MKS Pirani gauge and an MKS Mini-Ion were mounted onto each chamber to monitor chamber pressures. After the new equipment was mounted, the two-stage vacuum chamber was sealed and pumped down. After leak checks and minor adjustments, pressures less than $8 \times 10^{-9}$ Torr were achieved.
Adjustments were made to the interior of the chambers along the analyte flight path. A mounting plate was built and installed on the entry port to the first stage to allow a new Beam Dynamics skimmer cone to be placed in line with the quadrupole axis (see Figure 3.3). The skimmer currently installed has a 1.0 mm opening and is held in place by a metal collar; the vacuum seal is provided by a thin layer of TorrSeal between the skirt of the skimmer and the mounting plate. A second skimmer with an opening of 0.5 mm is available (but not yet mounted) if working pressures higher than 1 Torr are needed.

Figure 3.4 – A skimmer cone with a 1mm opening gives entry into the 1st stage.
Between the first and second stage, a 1 mm pinhole plate was designed to narrow the aperture between the stages. This plate also allows for a beam chopper to be mounted to interrupt the analyte path. (see Figure 4.5)

![Diagram of aperture plate](image)

Figure 3.5- The aperture plate that separates the 1st and 2nd chamber

Because the inherited vacuum chamber and the old Pyrex flow cell were originally part of two separate systems, a new intermediate chamber had to be designed to couple them together. From the new intermediate chamber extend five flanges. One radial flange with a KF connection is used to direct the main gas flow to the rotary pump in the channel. Four other
radial Conflat flanges are available for LIF detection (see Figure 3.6). A 125 cfm mechanical pump (Edwards Model E2M175) connected to the intermediate stage is employed to maintain a steady state pressure of about 1 Torr in the reactor.

The flow reactor cell consists of an 80-cm-long, 3.75 cm-i.d. Pyrex tube coated with halocarbon wax to reduce radical wall loss. Flows are controlled inside the reaction cell by MKS Flow controllers (2000 sccm, 200 sccm, 100 sccm, and 50 sccm). Pressure in the reaction cell is monitored using MKS Baratron 10 Torr and 1000 Torr gauges.

Figure 3.6 – An intermediate stage was machined to couple the vacuum chamber and glass flow cell.

A removable liquid nitrogen trap was placed downstream of the reactor in order to protect the vacuum pump from corrosive reactants and products. Several other small parts had to be
designed and machined in order to couple the vacuum chamber to the flow cell. For further interest, the supplementary material section contains most of the technical drawings used to machine necessary parts.

Initial leak testing uncovered minor leaks, that were then fixed. At the present time, with 1 Torr of He in the flow tube, the pressure in the 1st Stage chamber is about $1 \times 10^{-5}$ Torr and the 2nd Stage is at $3 \times 10^{-7}$. If no gas is flowing, the 2nd Stage has reached $3 \times 10^{-8}$ Torr. No bake-out or over-night pumping has been attempted yet. A calculated ratio of pressures between the flow cell and the 1st stage using a 1700 l/s through a 1 mm pinhole was compared to actual pressures observed in the system (see Figure 3.7).

![Figure 3.7](image_url)  
Figure 3.7 – Calculations were done for N$_2$, measurements with Ar.
The installed quadrupole mass spectrometer is an Extrel Model 150 Model 5221. Alignment of the quadrupole through the vacuum apertures was achieved using a red HeNe laser. Detection was carried out by continuous sampling at the downstream end of the flow tube, through a two-stage beam inlet system. The mass spectrometer was set to emit bombarding electrons with 25 eV of impact energy.

Initial testing of the ion signals was done using 1 Torr of He with a small amount of NO (Figure 3.8). Quadrupole settings were adjusted to optimize signal detection. At constant dynode voltage, the signal at mass 30 increased linearly with filament emission current between 1 and 4 milliamps (see Figure 3.9). At constant emission current, the signal increased rapidly as the dynode voltage was increased from 2.3 to 2.6 kilovolts. But voltages higher than 2.6 kV caused the signal to decrease (Figure 3.10). It was later found that the signal improved significantly when a 50 ohm-to-ground shunt was inserted at the input of the F-100T Amplifier/Discriminator.

![Signal from Approximately 2 mTorr NO in 0.98 He](image)

Figure 3.8 – Mass spectrum with emission current 1mA and dynode voltage 2.6kV
Figure 3.9 – Tuning the emission current on the quadrupole for optimal signal.

Figure 3.10 – Tuning the dynode voltage for optimal detection.
3.6 Detecting 2-Hydroxyethyl-Peroxy Radical

Attempts were made to make \( \cdot \text{CH}_2\text{CH}_2\text{OH} \) radical by reacting hydrogen atoms with 2-iodoethanol.

\[
\text{H} + \text{ICH}_2\text{CH}_2\text{OH} \rightarrow \text{HI} + \cdot \text{CH}_2\text{CH}_2\text{OH}
\]

The rate constant for this reaction has not been measured, but similar reactions of H with alkyl iodides are reasonably fast \( (k \sim 1 \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}) \). Hydrogen atoms were generated by flowing a mixture of 1% H\(_2\)/Argon through a microwave discharge. A separate flow of He was bubbled through liquid 2-iodoethanol at room temperature, on the low pressure side of the flow controller. The vapor pressure of iodoethanol is not known. The two flows were combined in the flow tube, approximately 70 cm from the pinhole. The sliding injector was removed for this experiment.

The current quadrupole setup uses ¾ inch diameter rods and a 2.1 MHz 300 watt power supply. According to the Extrel literature, this is designed for the mass range of 1 to 120 amu. The parent peak of iodoethanol is 172. While we have occasionally observed a signal at 172, its appearance depends on the quadrupole tuning and the peak is distorted and much weaker than it should be compared to lower mass peaks. Also for this experiment the masses at 127 (I atoms) and 128 (HI) are useful signals. These were consistently observed, but weaker than expected. The older quadrupole electronics that are in storage have a 1.2 MHz 100 watt power supply; according to Extrel this would extend the upper range to 200 amu.
In our system, the strongest peak observed for iodoethanol is at mass 45 (C$_2$H$_5$O$^+$), in agreement with the spectrum recorded in the NIST WebBook. (See Figure 3.11) According to the latter, the parent peak at 172 should be next strongest. Other significant peaks occur at 27, 31, 43 and 127.

When the hydrogen discharge is on, evidence for the above reaction is as follows: the signal at mass 45 decreases; the signals at 127 and 128 increase; at the conclusion of flowing for 4 hours, the cold trap protecting the rotary pump has a red deposit (iodoethanol is colorless); when the trap comes to room temperature, a strip of starch iodide paper turns blue, confirming the presence of I$_2$.

As the flow of He through the bubbler was decreased by a factor of two, and then by another factor of two, the amount of iodoethanol being destroyed, as judged by the decrease in the mass 45 signal, decreased by only $\sim 20\%$, suggesting that the limiting reagent was the hydrogen atom concentration, not the iodoethanol.
3.7 Conclusion

The DFMS has shown to be a viable working system. The chamber maintains high vacuum, allowing for low background in the mass spectra. The quadrupole mass spectrometer is functional and discriminating masses appropriately; however, its detection limits still need to be tested. The flow cell is holding vacuum and has no sign of leakage. All the parts, gauges, controllers, plumbing, electrical systems, computer software, vacuums, and monitors appear to be in full function allowing for a working discharge-flow mass spectrometer.

4 FUTURE WORK WITH DISCHARGE FLOW MASS SPECTROMETER

4.1 Summary

Future work with the DFMS will investigate the effects of water vapor on the kinetics and product branching ratio of the reaction of organic peroxy radicals with nitric oxide.

4.2 Peroxy Radicals in the Troposphere

Organic peroxy radicals (RO$_2$) in the atmosphere are intermediates in the combustion of hydrocarbons and serve as precursors for tropospheric ozone formation. Peroxy radicals are produced through hydrogen extraction by atmospheric radicals, followed by the addition of O$_2$ to the radical site on the hydrocarbon. The reaction mechanism is summarized by reactions 4.1 and 4.2:

$$ RH + (OH, Cl, NO_3) \rightarrow R + (H_2O, HCl, HNO_3) $$ (4.1)
\[ R + O_2 \rightarrow RO_2 \quad (4.2) \]

The reaction of \( RO_2 \) with NO produces two different products, as shown in reactions 4.3 and 4.4.

\[ RO_2 + NO \rightarrow RO + NO_2 \quad (4.3) \]

\[ RO_2 + NO \rightarrow RONO_2 \quad (4.4) \]

\[ NO_2 + h\nu \rightarrow NO + O \quad (4.5) \]

\[ O + O_2 \rightarrow O_3 \quad (4.6) \]

RO and RONO\(_2\) are intermediates to tropospheric pollution, and the photolysis of NO\(_2\) is the primary source of tropospheric ozone (Reactions 4.5 and 4.6). Alkoxy radicals (RO) react with oxygen to produce aldehydes and HO\(_2\). HO\(_2\) is the simplest and most abundant peroxy radical and contributes to pollution by the same mechanism as other peroxy radicals (reactions 4.3 and 4.4). Organic nitrates (RONO\(_2\)), the product of reaction 4.4, are stable and relatively inert. This inertness and stability permits organic nitrates to be carried by the wind to adjacent regions. They eventually break down (due to heat) into NO\(_2\) and RO radicals, which then allows for
reaction 4.3 to occur, leading to ozone production. This cascade of events provides a mechanism for pristine environments to become polluted.

Previous work has been done to understand the branching ratio of reactions 4.3 and 4.4 as a function of the R-group size, but no one has studied the kinetics or branching ratio of these reactions as a function of water-vapor. The branching ratio is defined as the ratio of reaction 4.3 and 4.4 \( \frac{k_{4.4}}{k_{4.4}} \).

Previous experimental work shows that the presence of water vapor perturbs reactions 4.5 and 4.6.

\[
\begin{align*}
\text{HO}_2 + \text{NO} &\rightarrow \text{OH} + \text{NO}_2 \quad (4.5) \\
\text{HO}_2 + \text{NO} &\rightarrow \text{HONO}_2 \quad (4.6)
\end{align*}
\]

The product branching ratio \( \frac{k_{4.5}}{k_{4.6}} \) is perturbed in favor of increased nitrate production (HONO\(_2\)). Butkovskaya et al.\(^7^4\) showed that at 50% relative humidity production of HNO\(_3\) increased by 800%. They hypothesized that water forms a mono-hydrated complex with peroxy radicals. An HO\(_2\)-H\(_2\)O complex has been observed experimentally by Suma et al.\(^3^4\) An analysis of published data by Butkovskaya et al. supports the hypothesis that water complexation of RO\(_2\) leads to increased nitrate production.\(^7^4\) Butkovskaya et al. depict the reported product branching ratio as a function of the percentage of HO\(_2\) that is complexed with water which shows a clear relationship between increased nitrate production and complexation with water.

Previous work\(^7^5\) demonstrates that organic peroxy radicals form water complexes and it is proposed that their reactions with NO can lead to increased production of organic nitrates in the troposphere. In a recent high level \textit{ab initio} study, Clark et al.\(^7^6\) demonstrated theoretically the
existence of other peroxy radical-water complexes. Perturbation to the product branching ratio favoring formation of RONO₂ is speculated to occur because of the formation of a radical-molecule complex. The production of NO₂ is caused by NO removing oxygen from the peroxy radical at the nitrogen site. The NO bond is formed when the NO₂ rotates to attach to RO. It is proposed that a water complex provides a chaperone whereby the NO₂ is maintained close to the RO through hydrogen bonding. The longer the NO₂ has to interact with the RO at close proximity, the more likely the formation of RONO₂. The work presented in Chapter 2 with the β-HEP self-reaction supported the theory that complexation of peroxy radicals with water increases the rate of reaction. Although this discovery is of importance to the atmospheric science community, the perturbation of the branching ratio of β-HEP + NO towards more organic nitrate formation has more direct effects on the atmosphere due to the ability to pollute pristine environments through long-range transport.

4.3 Kinetic Model of HydroxyethylPeroxy Radical and NO

A kinetic model of the reaction mechanism was built using FACSIMILE to aid in the analysis of the kinetic data. This allows for secondary chemistry to be explored in order to assure that significant sources, sinks, or pathways of analyte molecules are accounted for in rate calculations.

\[
\text{HOCH}_2\text{CH}_2\text{O}_2 + \text{wall} \rightarrow \text{Products} \quad (4.7)
\]

\[
\text{HOCH}_2\text{CH}_2 + \text{O}_2 \rightarrow \text{HOCH}_2\text{CH}_2\text{O}_2 \quad (4.8)
\]

\[
\text{HOCH}_2\text{CH}_2\text{O}_2 + \text{NO} \rightarrow \text{HOCH}_2\text{CH}_2\text{O} + \text{NO}_2 \quad (4.9)
\]
The reactions outlined above constitute the reactions used in the model. This model allows data to be fit and analyzed for determination of the rate constant, and also aids in predicting probable effects from water vapor. These reactions are a small sample to demonstrate the complexity inside the cell. To minimize effects from undesired chemistry, the concentration of ICH₂CH₂OH in the model is orders of magnitude smaller than [O₂]. This will produce pseudo first-order conditions whereby the relative concentrations of O₂ will remain unchanged. Figure 4.1 show the results of the model run for 40 milliseconds. As suspected, CH₂CH₂OH reacts
quickly with O₂ to form HEP. As HEP reacts with NO, both NO₂ and HOCH₂CH₂ONO₂ are formed.

**4.4 Experimental Method**

The detection method for the HEP + NO reaction must include sufficient selectivity, sensitivity, and response time to obtain the necessary kinetic information. It is necessary to measure the time-resolved concentration of the reactants and products simultaneously. The DFMS described in Chapter 3 should be suitable for this reaction.

Figure 4.2 outlines the proposed settings of the DFMS. The primary He gas flow will introduce various water-vapor concentrations into the system. Atomic hydrogen will be produced through a microwave discharge and introduced into the cavity. Through the sliding injector,
HOCH₂CH₂I along with O₂ will be added to the system. The appropriate NO, NO₂, CH₂CH₂OH, HOCH₂CH₂OO, HI peaks will monitored on the mass spectra

Figure 4.2 – DFMS setup for studying the kinetics of the HEP + NO reaction
5 SUMMARY

This thesis was written to present studies related to the field of atmospheric chemistry. The first portion outlines the importance for understanding atmospheric reactions and their relation to pollution. The $\beta$-HEP self-reaction results, combined with the work from Butkovskaya on water enhancement of HO$_2$ + NO, provides motivation for further investigation into the effects of water on other peroxy radical reactions. To further these investigations, a discharge flow mass spectrometer was engineered. The DFMS serves as a simple and ideal way to investigate water vapors effects on peroxy radical reactions. To correctly predict pollution levels, atmospheric models need to be adjusted to include the rate enhancements demonstrated in this work. The implications of these findings promote the path of further discovery of peroxy radical reactions enhancements through the formation of pre-reactive complexes between radicals and polar molecules such as water, methanol, and ammonia. Furthermore, water vapor catalysis will become increasingly more significant in light of global warming increasing water vapor concentrations in the troposphere.
BIBLIOGRAPHY


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SUPPLEMENTARY

Technical Drawings

The technical drawings include the technical drawings of the parts designed and machined for the discharge flow mass spectrometer. The machining was done by Hamilton Tools and Engineering, MDC Vacuum, or 80/20 Inc.
This is the collar that clamps down over the skimmer cone onto the skimmer mounting plate.
This is mounting plate for the skimmer cone. This covers the entrance into the first stage of the vacuum chamber.
This is the collar that goes over that glass flow cell. It clamps the glass cell to the new intermediate stage, providing the pressure needed to create the seal around the O-ring.
A cut view of the glass flow cell collar.
Dimensions of the glass cell flow tube.
A sketch of how the vacuum chamber (A), the skimmer cone plate (B), the intermediate stage (C), and the flow cell (D) couple together.
Measurements of the vacuum chamber coupled to the flow cell via the intermediate stage.
Dimensions of the intermediate stage.
Dimensions of the intermediate stage.
Dimensions of the intermediate stage. This face mounts with the vacuum chamber.
Dimensions of the intermediate stage. This face mounts with the glass cell.
Dimensions of the intermediate stage.
Intermediate stage mounted to the vacuum cage.
Flanges that cover the LIF ports on the intermediate stage.
Support beams for the flow cell. These are adjustable up-and-down and side-to-side to allow optimal positioning of the flow cell into the intermediate stage.
Mounting plates for the flow cell support beams to the bread board table top.
The valve that seals the channel vacuum pump to the intermediate stage.
These flanges mount to a port on the vacuum chamber and on the intermediate chamber. They allow the vacuum chamber to be brought back up to atmospheric pressure.
Designs for a table. This table supports the flow cell and couples the vacuum chamber support to another table top.
S-1 Computational Methods

S.1.1 Partition function calculations

From the Gaussian 03 calculations, the partition function of each conformation of each molecule was calculated according to the theory by McQuarrie and Simon. The overall partition function for a molecule is approximated by the product of its translational, rotational, electronic, and vibrational partition functions, which are assumed to be separable. The expressions for the translational, rotational, and electronic partition functions are:

\[ q^T = \frac{V (2\pi mkT)^{3/2}}{h^3} \]  

(S1)

where \( V \) is the volume of the reaction cell, \( m \) is the mass of the molecule, \( k \) is Boltzmann's constant, \( h \) is Planck's constant, and \( T \) is temperature;
\[ q^R = \left( \frac{kT}{hc} \right)^{3/2} \left( \frac{\pi}{ABC} \right)^{1/2} \] (S2)

where \(c\) is the speed of light and \(A, B,\) and \(C\) are the rotational constants of the molecule;

\[ q^E = ge^{-E/(RT)} \] (S3)

where \(g\) is the degeneracy of the electronic ground state and \(E\) is the zero-point energy of the ground vibrational state. Note that the energy in the partition function for all molecules and geometries must be in reference to the same reference energy.

**S.1.2 Vibrational partition function**

The vibrational partition function for a molecule is the product of the partition functions for each of its vibrational modes, assuming that the normal vibrational modes of the molecule are independent. The partition function of a vibrational mode can be calculated as the sum of the contributions from each vibrational state,

\[ q^V = \sum_v e^{-\beta E_v} \] (S4)
where $\beta = (kT)^{-1}$. The energies of the vibrational states can be calculated according to several models for the vibrational motions. In this calculation, the harmonic oscillator, Morse oscillator, and hindered rotor models were used for each mode according to the model that best approximated the vibrational motion of the mode.

**S.1.3 Harmonic oscillator**

A harmonic approximation assumes that the energy levels of a vibrational mode are equally spaced. This approximation is accurate for the lowest vibrational states and therefore can be made when only the ground and first excited states are occupied. For a harmonic oscillator, equation S4 becomes

$$ q^V = \sum_{\nu} e^{-\beta \hbar c \tilde{\nu}} = \frac{1}{1 - e^{-\beta \hbar c \tilde{\nu}}} $$

(S5)

where $\tilde{\nu}$ is the fundamental frequency of the vibrational mode. In this calculation, the fundamental anharmonic frequency calculated in Gaussian was used for $\tilde{\nu}$. 

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A Morse oscillator can be used to model vibrational modes that are dissociative. A harmonic oscillator model assumes that all vibrational states are equally spaced and does not account for the possibility that a bond can dissociate with sufficient energy. Therefore, the partition function based on the harmonic oscillator tends to underestimate the true partition function of the mode. The Morse oscillator accounts for the decreasing spacing between the vibrationally excited states and eventually the dissociation of the bond.

The energy levels for the Morse oscillator potential are given by,

\[ G(v) = v_e \left( v + \frac{1}{2} \right) + x_e \left( v + \frac{1}{2} \right)^2 \]  

(S6)

where \( v_e \) is the fundamental harmonic vibrational frequency, \( v \) is the vibrational quantum number, and \( x_e \) is the diagonal element of the \( X \)-matrix corresponding to the vibrational mode. If \( x_e \) is negative, the bond will eventually dissociate, whereas if \( x_e \) is positive there are infinitely many bound states. The bond will dissociate when \( G(v) \) achieves a maximum, or at \( v^* = -\frac{v_e}{2x_e} - \frac{1}{2} \).

Therefore, at the highest energy bound state, \( v_{max} = \left\lfloor v^* \right\rfloor \) and there are \( N = v_{max} + 1 \) bound states. For some states, if the vibration primarily involved the stretching of a hydrogen bond, \( \hat{x}_e \) was calculated from the dissociation energy for the breaking of the hydrogen bond, \( D \),
\[ \tilde{X}_e = -\frac{V_e^2}{4D} \]  

(S7)

\( N \) was then calculated as \( \frac{V_e}{2\tilde{X}_e} \) and rounded to the nearest integer. The partition function was then calculated,

\[ q^V = \sum_{v=0}^{v_{max}} e^{-\frac{\hbar c (G(v) - G(0))}{kT}} \]  

(S8)

from \( \tilde{X}_e \) to calculate \( G(v) \) and \( G(0) \).

### S.1.5 Hindered rotor

A hindered rotor model was used to model vibrations that involve the rotation of a functional group on the molecule. The calculation of the partition function for a hindered rotor vibration was based on McClurg et al. \[^{42}\]
\begin{equation}
q' = q_{HO} \left( \frac{r \pi}{\theta} \right)^{1/2} e^{-r(2\theta)} I_0 \left( \frac{r}{2\theta} \right) \tag{S9}
\end{equation}

where \( q_{HO} \) is the partition function calculated as a harmonic oscillator using the fundamental harmonic frequency, and \( r \) and \( \theta \) are defined as

\begin{equation}
r = \frac{\sqrt{2Iw}}{\hbar} \tag{S10}
\end{equation}

\begin{equation}
\theta = \frac{kT}{\hbar} \sqrt{\frac{2I}{w}} \tag{S11}
\end{equation}

where \( w \) is the barrier height for the hindered rotor (for this complex, the strength of one or two hydrogen bonds), and \( I \) is the moment of inertia for the rotation.

\subsection*{S.1.6 Local minima weighting}

Three local minimum geometries for the complex and two local minimum geometries for \( \beta \)-HEP are all accessible at room temperature. Therefore the partition function for each of these
geometries contributes to the overall equilibrium constant for complex formation. The partition function for a molecule is

\[ q = \sum_i e^{-\beta E_i / (kT)} \]  

where \( i \) denotes all of the states for the molecule. Therefore the partition functions for \( \beta \)-HEP-H\(_2\)O and \( \beta \)-HEP are equal to the sums of the partition functions for each of the local minimum geometries if they have a common reference energy. Therefore, the equilibrium constant for the complex formation is equal to

\[
K_e = \frac{\frac{q_{(\text{HEP-H}_2\text{O})_1} + q_{(\text{HEP-H}_2\text{O})_2} + q_{(\text{HEP-H}_2\text{O})_3}}{V}}{\left(\frac{q_{(\text{HEP})_1} + q_{(\text{HEP})_2}}{V}\right) \times \left(\frac{q_{\text{H}_2\text{O}}}{V}\right)}
\]  

(S13)
Figure S 1: Potential energy surface for Cl + HOCH\textsubscript{2}CH\textsubscript{2}Cl reaction

Figure S-1: Potential energy surface diagram for the reaction of Cl + HOCH\textsubscript{2}CH\textsubscript{2}Cl calculated at MP2/6-311++G(3df,3pd). Energies reported in kcal mol\textsuperscript{-1}. The energies of both reactants (Cl and HOCH\textsubscript{2}CH\textsubscript{2}Cl), pre-reactive complex and transition states were optimized using the B3LYP/6-311++G(3df,3pd) method/basis set. Single point calculations were computed at the

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MP2/6-311++G(3df,3pd) level for all species using the optimized geometry computed at the B3LYP/6-311++G(3df,3pd) level.

S-2 Calculation of concentrations of O₃, HO₂, β-HEP and β-HEP-H₂O from measured absorbances

From the Beer-Lambert Law, where \( b \) is pathlength, \( \sigma \) is absorption cross section, and \( [\cdot] \) indicates concentrations

\[
A_{220} = b \left( \sigma_{220}^{O_3}[O_3] + \sigma_{220}^{HO_2}[HO_2] + \sigma_{220}^{\beta-HEP}[\beta - HEP] + \sigma_{220}^{\beta-HEP-H_2O}[\beta - HEP - H_2O] \right)
\]

(S14)

\[
A_{230} = b \left( \sigma_{230}^{O_3}[O_3] + \sigma_{230}^{HO_2}[HO_2] + \sigma_{230}^{\beta-HEP}[\beta - HEP] + \sigma_{230}^{\beta-HEP-H_2O}[\beta - HEP - H_2O] \right)
\]

(S15)
\[ A_{254} = b \left( \sigma_{254}^{O_3}[O_3] + \sigma_{254}^{HO_2}[HO_2] + \sigma_{254}^{\beta - HEP}[\beta - HEP] + \sigma_{254}^{\beta - HEP-H_2O}[\beta - HEP - H_2O] \right) \]  

(S16)

Since the cross sections of the \( \beta \)-HEP-H\(_2\)O complex is unknown, we assume

\[ \sigma_x^{\beta - HEP - H_2O} = f_x \sigma_x^{\beta - HEP} \]  

(S17)

Where \( f_x \geq 0 \) and \( x = 200 - 400 \) nm.

Substituting Equation S17 into the last two terms of equation S14,15 and 16 yields,

\[ \sigma_x^{\beta - HEP}[\beta - HEP] + \sigma_x^{\beta - HEP-H_2O}[\beta - HEP - H_2O] = \sigma_x^{\beta - HEP}[\beta - HEP] + f_x \sigma_x^{\beta - HEP}[\beta - HEP - H_2O] \]  

(S18)

The equilibrium constant for the formation of the \( \beta \)-HEP-H\(_2\)O complex is,
\[ K = \frac{[\beta - \text{HEP} - H_2O]}{[\beta - \text{HEP}][H_2O]} \]

Solving for \( \beta \)-HEP-H_2O gives,

\[ [\beta - \text{HEP} - H_2O] = K[\beta - \text{HEP}][H_2O] \]  
(S19)

Substituting Equation S19 into S18 yields,

\[
\sigma_x^{\beta-\text{HEP}}[\beta - \text{HEP}] + \sigma_x^{\beta-\text{HEP}-H_2O}[\beta - \text{HEP} - H_2O] = \sigma_x^{\beta-\text{HEP}}[\beta - \text{HEP}] + f_x K \sigma_x^{\beta-\text{HEP}}[\beta - \text{HEP}][H_2O] 
\]
(S20)

Factoring out \( \sigma_x^{\beta-\text{HEP}}[\beta - \text{HEP}] \) yields Equation S21

\[
\sigma_x^{\beta-\text{HEP}}[\beta - \text{HEP}] + \sigma_x^{\beta-\text{HEP}-H_2O}[\beta - \text{HEP} - H_2O] = \sigma_x^{\beta-\text{HEP}}[\beta - \text{HEP}][1 + f_x K[H_2O]] 
\]
(S21)

Thus the equations for absorbance at 220, 230 and 254 reduce to,
\[ A_{220} = b\left(\sigma_{220}^{O_3}[O_3] + \sigma_{220}^{HO_2}[HO_2] + \sigma_{220}^{\beta-HEP}[\beta-HEP]\{1 + f_{230}K[H_2O]\}\right) \]

(S22)

\[ A_{230} = b\left(\sigma_{220}^{O_3}[O_3] + \sigma_{220}^{HO_2}[HO_2] + \sigma_{220}^{\beta-HEP}[\beta-HEP]\{1 + f_{230}K[H_2O]\}\right) \]

(S23)

\[ A_{254} = b\left(\sigma_{220}^{O_3}[O_3] + \sigma_{220}^{HO_2}[HO_2] + \sigma_{220}^{\beta-HEP}[\beta-HEP]\{1 + f_{254}K[H_2O]\}\right) \]

(S24)

If, \(0 \leq f_{220}, f_{230}, f_{254} \leq 10\), the product of \(f_{220}K[H_2O], f_{230}K[H_2O], f_{254}K[H_2O]\) \(\approx 0\). Thus the observed \(\beta\)-HEP concentration in this experiment is,

\[ [\beta-HEP]_{\text{observed}} = [\beta-HEP]\{1 + f_xK[H_2O]\} \approx [\beta-HEP] \]

(S25)

If \(f_x=0\), the two cross sections do not overlap and the observed signal in this experiment is only due to \(\beta\)-HEP. If \(f_x > 0\) there is an overlap between the two cross sections and the observed
signal in this experiment is due to both β-HEP and β-HEP-H₂O complex. In either case, the equations to solve for the concentrations of ozone, β-HEP and HO₂ reduce to

\[
A_{220} = b\left(\sigma_{220}^{O_3}[O_3] + \sigma_{220}^{HO_2}[HO_2] + \sigma_{220}^{\beta-HEP}[^\beta-HEP]\right)
\]

(S26)

\[
A_{230} = b\left(\sigma_{230}^{O_3}[O_3] + \sigma_{230}^{HO_2}[HO_2] + \sigma_{230}^{\beta-HEP}[^\beta-HEP]\right)
\]

(S27)

\[
A_{254} = b\left(\sigma_{254}^{O_3}[O_3] + \sigma_{254}^{HO_2}[HO_2] + \sigma_{254}^{\beta-HEP}[^\beta-HEP]\right)
\]

(S28)

The HO₂ concentration calculated by solving equation S26, 27, 28 was always below our detection limit of 1×10¹³ molecules/cm³. Thus the general form of the measured absorbance is given by Equation 1 in the paper

\[
A_x = b\left(\sigma_x^{O_3}[O_3] + \sigma_x^{\beta-HEP}[^\beta-HEP]\right)
\]

Equation 1
S-3 Kinetic Equations

To derive Equation 2, we define the following variables

Concentration of un-complexed $[\beta\text{-HEP}] = [u]$

Concentration of complexed $\beta\text{-HEP} [\beta\text{-HEP-H}_2\text{O}] = [c]$

Concentration of observed $\beta\text{-HEP} = [o]$

Concentration of water $= [w]$

Observed rate constant $= k_o$

From elementary reactions 3 and 5, the rate of loss of un-complexed $\beta\text{-HEP}$ is given by

$$\frac{d[u]}{dt} = -2k_3[u]^2 - k_5[u][c]$$

(S29)

From elementary reactions 5 and 6, the rate of loss of $\beta\text{-HEP-H}_2\text{O}$ is given by

$$\frac{d[c]}{dt} = -k_5[u][c] - 2k_6[c]^2$$

(S30)
To obtain \( k_3, k_5 \) and \( k_6 \) as functions of \( k_o \) we introduce the relationship,

\[
\frac{d[o]}{dt} = -2k_o[o]^2
\]

(S31)

Substituting equation S25 (the observed concentration of \( \beta \)-HEP in our experiment) in equation S31 yields,

\[
\frac{d[o]}{dt} = -2k_o([u][1 + fK[w]])^2
\]

\[= -2k_o[u]^2([1 + fK[w]])^2\]

(S32)

Adding Equations S-29 and S-30 gives the theoretically observed loss of \( \beta \)-HEP in this experiment

\[
\frac{d[o]}{dt} = \frac{d[u]}{dt} + \frac{d[c]}{dt} = -2k_3[u]^2 - 2k_5[u][c] - 2k_6[c]^2\]

(S33)
Equating S-32 and S-33 gives,

\[-2k_o[u]^2([1 + fK[w]])^2 = -2k_3[u]^2 - 2k_5[u][c] - 2k_6[c]^2\]

(S34)

The equilibrium constant for complex formation is given by

\[K = \frac{[c]}{[u][w]}\]

Thus,

\[[c] = K[u][w]\]

(S35)

Substituting Equation S-35 into Equation S34 gives,

\[-2k_o[u]^2([1 + fK[w]])^2 = -2k_3[u]^2 - 2k_5[u]^2K[w] - 2k_6[u]^2[w]^2K^2\]

(S36)

Dividing Equation S36 by -2[u]^2 on both sides yields,
\[ k_o([1 + fK[w]]^2 = k_3 + k_5[u]K[w] + k_6[u]^2[w]^2K^2 \]

Solving for \( k_o \) gives,

\[ k_o = \frac{k_3 + k_5[u]K[w] + k_6[u]^2[w]^2K^2}{([1 + fK[w]]^2} \]

(S37)

From calculated values of the equilibrium constant (Figure 8) and the slopes of the experimental data (Figure 7), \( k_6[u]^2[w]^2K^2 << k_3 + k_5[u]K[w] \). Further within the uncertainty, there is no evidence for curvature in data in Figure 7.

Also, if \( f = 1 \) the cross section of \( \beta \)-HEP and the complex are equal then \( 1 + fK[w] \approx 1 \). If \( f = 0 \) the cross section of \( \beta \)-HEP and the complex are distinguishable, then \( 1 + fK[w] = 1 \). In either case Equation S -37 reduces to Equation 2 in the manuscript as,

\[ k_o = k_3 + k_5K[w] \]

Equation 2
**S-4 Derivation of Equation 3**

Following is the derivation of equation of Equation 3 from Equation 2

\[ k_o = k_3 \left\{ 1 + \frac{k_5}{k_3} K [w] \right\} \]  

(S38)

From the Arrhenius expression for the temperature dependence of the rate coefficients and van’t Hoff equation for the temperature dependence of the equilibrium constant,

\[ k_3 = A_3 e^{-\frac{E_3}{RT}} \]  

(S39a)

\[ k_5 = A_5 e^{-\frac{E_5}{RT}} \]  

(S39b)

\[ K = C * e^{-\frac{\Delta H}{RT}} \]  

(S39c)
Substituting S-39a, S39b, S39c into Equation S-38 gives,

\[ k_o = k_3 \left\{ 1 + \frac{A_5 e^{\frac{-E_5}{RT}} C e^{\frac{-\Delta H}{RT}} [w]}{A_3 e^{\frac{-E_3}{RT}}} \right\} \]

\[ k_o = k_3 \left\{ 1 + \frac{A_5 * C}{A_3} e^{\frac{-(\Delta H + E_5 - E_3)}{RT}} [w] \right\} \]

(S40)

Combining parameters, \( A_5 C / A_3 = A \) and \( (E_5 + \Delta H - E_3) = E \) reduces Equation S40 to Equation 3 in the manuscript.

\[ k_o = k_3 \left\{ 1 + A e^{\frac{-E}{RT}} [w] \right\} \]

Equation 3