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Time-Varying MIMO Channels: Parametric Statistical Modeling and Experimental Results

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Abstract—Accurate characterization of multiple-input multiple-output (MIMO) fading channels is an important prerequisite for the design of multiantenna wireless-communication systems. In this paper, a single-bounce two-ring statistical model for the time-varying MIMO flat fading channels is proposed. In the model, both the base station and mobile station are surrounded by their own ring of scatterers. For the proposed model, a closed-form expression for the spatio-temporal cross-correlation function between any two subchannels is derived, assuming single-bounce scattering. The new analytical expression includes several key physical parameters of interest such as the mean angle-of-departure, the associated angle spreads, and the Doppler spread in a compact form. The model includes many existing correlation models as special cases. Its utility is demonstrated by a comparison with collected MIMO data in terms of the spatio-temporal correlations, level crossing rate, average fade duration, and the instantaneous mutual information.

Index Terms—Average fade duration (AFD), channel modeling, instantaneous mutual information (IMI), level crossing rate (LCR), multiantenna systems, multiple input multiple output (MIMO), Rayleigh fading, spatio-temporal cross correlation (STCC).

I. INTRODUCTION

THE UTILIZATION of antenna arrays at the base station (BS) and the mobile station (MS) in a wireless-communication system increases the capacity linearly with $\min(N_T, N_R)$, where $N_T$ and $N_R$ are the numbers of transmit (Tx) and receive (Rx) antenna elements, respectively, provided that the environment is sufficiently rich in multipath components [3], [4]. The significant increase of channel capacity was originally reported for uncorrelated subchannels [5], [4]. However, later, it was realized that the correlation among subchannels can significantly affect capacity [5]–[8] and other system performance metrics [9]–[12]. For an overview of multiple-input multiple-output (MIMO) channel models, see [13]–[16] and the references therein. In this paper, we concentrate on the statistical modeling of time-varying MIMO flat fading channels and verify the utility of the proposed model via comparison with different sets of data.

In a typical macrocell, the BS is elevated and receives the signal within a narrow beamwidth, whereas the MS is surrounded by local scatterers. MIMO modeling of this typical macrocell environment is investigated in [5], where a closed-form expression for the MIMO spatio-temporal cross correlation (STCC) is derived, assuming nonisotropic scattering around the BS. However, in the outdoor and indoor picocells, both the BS and MS are normally surrounded by local scatterers. Clearly, the MIMO macrocell model of [5] cannot be used for such environments. For this situation, we need a double-directional channel model (see [17] and [18], where the double-directional concept is introduced, and some measurements results are provided).

In this paper, the space-time model of [5] is extended by adding another ring of scatterers around the BS in order to facilitate a new and mathematically tractable MIMO STCC model that is needed for analytical calculations and system design, for example, array optimization [19]. We call the newly extended model as the single-bounce two-ring model since only single-bounce rays are considered. The proposed model includes [5] and some other existing correlation models as special cases. Note that the proposed model belongs to the class of double-directional channel models since it includes angular information at both the BS and MS. Most importantly, the proposed parametric model provides a compact analytical form for MIMO STCC in terms of several key physical parameters of the channel. It also avoids the technical difficulties of the double-bounce two-ring model (e.g., [20]), which is discussed in [13]. The parametric nature of the model makes it adaptable to a variety of propagation environments, and its

1In this paper, each subchannel represents the radio link between each Tx/Rx pair of antennas with the time-varying gain in (2).

2Only the flat fading is considered here since the MIMO frequency-selective fading channel can be converted into several flat fading channels by the orthogonal frequency-division-multiplexing (OFDM) technique.
compact mathematical form is convenient for both analytical studies and numerical calculations. Finally, comparison of the model with the data collected at Helsinki University of Technology (HUT) [21], [22] and Brigham Young University (BYU) [23], [24] confirms the utility of the model in real environments in terms of mutual information, several types of correlations, level crossing rate (LCR), and average fade duration (AFD).

The rest of this paper is organized as follows. The single-bounce two-ring model is presented in Section II. The closed-form expression for the STCC function between any two subchannels is derived in Section III. Section IV includes data processing and parameter-estimation techniques. In Section V, a useful approximation when compared with the measured data. Finally, concluding remarks are provided in Section VI.

Notation: $\| \cdot \|_F$ is reserved for the Frobenius norm, $(\cdot)^*$ for the complex conjugate, $\mathbf{I}$ for the identity matrix, $(\cdot)^T$ for the transpose, and $\sqrt{-1}$ for the mathematical expectation, $j$ for $\sqrt{-1}$. $\mathbf{I}_m$ denotes the $m \times m$ identity matrix, and $[\mathbf{A}]_{m,n}$ is the $(m,n)$th element of the matrix $\mathbf{A}$. $t \in [m,n]$ with integers $m$ and $n$ implies that $t$ is an integer such that $m \leq t \leq n$, whereas $t \in [z_1,z_2]$ with real $z_1$ and $z_2$ denotes that $t$ is real such that $z_1 \leq t < z_2$. $\min(x,y)$ and $\max(x,y)$ indicate the minimum and maximum of real $x$ and $y$, respectively. vec($\mathbf{A}$) stacks all the columns of the matrix $\mathbf{A}$ into one tall column vector.

II. PROPOSED STATISTICAL MODEL

For a linear time-varying flat fading $N_R \times N_T$ MIMO system, the input–output relationship can be written as

$$r(t) = \sqrt{\frac{P_{\text{Tx}}}{N_T}} \mathbf{H}(t) \mathbf{s}(t) + e(t).$$

(1)

In (1) $P_{\text{Tx}}$ is the total emitted power from the Tx antenna array. $[\mathbf{H}(t)]_{l,p} = h_{lp}(t)$, $h_{lp}(t)$ denotes the complex baseband equivalent channel gain between the $p$th transmitter and the $l$th receiver. $r(t) = [r_1(t), r_2(t), \ldots, r_{N_T}(t)]^T$ is the received signal vector at time $t$. $\mathbf{s}(t) = [s_1(t), s_2(t), \ldots, s_{N_T}(t)]^T$ is the transmitted symbol vector from $N_T$ Tx antennas at time $t$ such that each element of $\mathbf{s}(t)$ has a unit power. $e(t) = [e_1(t), e_2(t), \ldots, e_{N_R}(t)]^T$ is the additive white Gaussian noise (AWGN) vector with the covariance matrix $\mathbf{R} = \mathbb{E}[e(t)e^H(t)] = P_{\text{Noise}} \mathbf{I}_{N_R}$ [5], where $P_{\text{Noise}}$ is the average noise power at each Rx antenna element.

The geometry of the proposed model is shown in Fig. 1 for a $2 \times 2$ channel where scatterers that are local to the BS and MS are modeled to be distributed on two separate rings. The key difference between our model and the other two-ring model [13] is that only single-bounce rays are considered, and multiple bounces are treated as secondary effects. This avoids the problem of the double-bounce two-ring model discussed in [13], facilitates the derivation of closed-form results, and proves to be a useful approximation when compared with the measured data.

In Fig. 1, which shows the forward channel (from the BS to the MS), the MS receives single-bounce rays from the scatterer $S_i$ around the MS (shown by dotted lines) and the scatterer $S'_k$ around the BS (shown by dash-dotted lines). $D$ is the distance between the MS and BS, $R$ and $R'$ are the radii of the scattering rings around the MS and BS, respectively, and $\alpha_{pq}$ and $\beta_{lm}$ are the directions of the subarray BS$_p$-BS$_q$ with element spacing $\delta_{pq}$ and subarray MS$_l$-MS$_m$ with element spacing $d_{lm}$, respectively.

For the frequency flat subchannel between the antenna elements BS$_p$ and MS$_l$, $h_{lp}(t)$ denotes the time-varying complex baseband equivalent channel gain. Mathematical representation of the superposition of rays at the MS results in the following expression for the channel gain:

$$h_{lp}(t) = \sqrt{\eta_{\text{Omega}} \sum_{k=1}^{N_R} \frac{g_k}{\sqrt{N'}} \exp\left[\frac{j\psi_k - 2\pi}{\lambda} (\xi_{pk} + \xi_{kl})\right]} + j2\pi f_D \cos(\varphi_k - \gamma) t + \sqrt{\eta_{\text{Omega}} \sum_{i=1}^{N} \frac{g_i}{\sqrt{N'}} \exp\left[\frac{j\psi_i - 2\pi}{\lambda} (\xi_{pi} + \xi_{il})\right]} + j2\pi f_D \cos(\phi_i - \gamma) t$$

(2)

Note that $\alpha_{pq} = \alpha_{pq} + \pi$, $p < q$, $\beta_{lm} = \beta_{lm} + \pi$, $l < m$, $\delta_{pq} = \delta_{pq} + \pi$, and $d_{lm} = d_{lm} + \pi$, with $p, q, l, m$ being the positive integers.
where the first and second summations correspond to the BS and MS rings, respectively. From (2), it is clear that the angle of arrival (AoA) and the angle of departure (AoD) play the interrelation between the single-bounce two-ring model in Fig. 2 and the channel transfer function $\mathbf{H}(t)$ in (1).

In (2), $\Omega_{lp}$ is the power transmitted through the subchannel BS$_p$-MS$_l$, i.e., $\Omega_{lp} = \mathbb{E}[|h_{lp}|^2]$, $\eta'$ and $\eta$ show the respective contributions of scatterers around the BS and MS to $\Omega_{lp}$ such that $\eta' + \eta = 1$, and $N$ and $N'$ are the numbers of scatterers around the MS and BS, respectively. The positive random variables $g_i$ and $g'_k$ represent the amplitudes of the waves scattered from $S_i$ and $S'_k$, which also include the antenna gains at $\phi_i$ and $\phi'_k$, and $\psi_i$ and $\psi'_k$ are the associated phase shifts. Furthermore, as shown in Fig. 2, $\phi'_k$ and $\phi_i$ are the AoDs of the waves that impinge on $S'_k$ and $S_i$, whereas $\phi_k$ and $\phi_i$ are the AoAs of the waves scattered from $S'_k$ and $S_i$. $\xi_{pk}$ and $\xi_{km}'$, which are functions of $\phi'_k$, are the lengths of BS$_p$-$S'_k$ and BS$_p$-$S_k'$ links, respectively, whereas $\xi'_{pq}$ and $\xi'_{km}$, which are functions of $\phi'$, are the lengths of $S'_k$-$MS_l$ and $S'_k$-$MS_m$ links. Other $\xi$'s can be easily identified in Fig. 2, and their functional relationships are given in (26a)–(26h). Finally, $\lambda$, $f_D$, and $\gamma$ are the wavelength, maximum Doppler frequency, and the direction of the MS motion, respectively.

The sets $\{g_i\}_{i=1}^N$ and $\{g'_k\}_{k=1}^{N'}$ consist of independent positive random variables with finite variances, independent of $\{\psi_i\}_{i=1}^N$ and $\{\psi'_k\}_{k=1}^{N'}$. We assume that $\{\psi_i\}_{i=1}^N$ and $\{\psi'_k\}_{k=1}^{N'}$ are the independent identically distributed (i.i.d.) random variables with uniform distributions over $[0, 2\pi)$. We also set $N^{-1}\sum_{i=1}^N g_i^2 = 1$ and $N'^{-1}\sum_{k=1}^{N'} g'_k^2 = 1$, which result in the desired identity $\mathbb{E}[|h_{lp}(t)|^2] = \Omega_{lp}$. The second moments of $g_i$, $g'_k$, and $g'_k$, $\forall k$ will be discussed in Section III. According to Fig. 2, $\phi_i$ is a function of $\phi_i$, and $\phi'_k$ is a function of $\phi'_k$; therefore, only $\phi'_k$ and $\phi_i$ are the independent angular variables. In what follows, we call $\phi'_k$ the AoD and $\phi_i$ the AoA.

III. MIMO Spatio-Temporal Cross Correlation

Based on the independent properties of $g_i$, $g'_k$, $\psi_i$, and $\psi'_k$, the normalized STCC between the two subchannel gains $h_{lp}(t)$ and $h_{mq}(t)$, which are defined as $\rho_{lp,mq}(\tau) = \mathbb{E}[h_{lp}(t + \tau)h_{mq}(t)]/\sqrt{\Omega_{lp}\Omega_{mq}}$, can be asymptotically written as

$$
\rho_{lp,mq}(\tau) = \lim_{N' \to \infty} \left(1 - \frac{\eta}{N'}\right) \sum_{k=1}^{N'} \mathbb{E}[g'_k^2] \times \exp\left[-\frac{2\pi}{\lambda} (\xi'_{pk} - \xi'_{qh} + \xi'_{kl} - \xi_{km})\right] + j2\pi f_D \cos(\phi'_k - \gamma) + \lim_{N' \to \infty} \frac{\eta}{N'} \times \sum_{i=1}^{N} \mathbb{E}[g^2_i] \exp\left[-\frac{2\pi}{\lambda} (\xi_{pi} - \xi_{q} + \xi_{il} - \xi_{im})\right] + j2\pi f_D \cos(\phi_i - \gamma)ight]. \quad (3)
$$

For large $N'$ and $N$, the small amounts of power received from $S_l$ and $S_i$ are proportional to $N'^{-1}\mathbb{E}[g'_k^2]$ and $N^{-1}\mathbb{E}[g_i^2]$, respectively. They are also equal to the infinitesimal powers from $\phi_i$ and $\phi'_k$ with probabilities $f_{\text{BS}}(\phi_i)d\phi_i$ and $f_{\text{MS}}(\phi'_k)d\phi'_k$, respectively. This implies that $N'^{-1}\mathbb{E}[g'_k^2] = f_{\text{BS}}(\phi'_k)d\phi'_k$ and $N^{-1}\mathbb{E}[g_i^2] = f_{\text{MS}}(\phi_i)d\phi_i$, where $f_{\text{BS}}(\phi)$ is the probability density functions (pdf) of the AoD, and $f_{\text{MS}}(\phi)$ is the pdf of the AoA. To simplify the notation, we define $x = \phi'_k$, $y = \phi_i$, $v = \phi'_k$, and $w = \phi_i$. Therefore, (3) can be reduced to the following integral form:

$$
\rho_{lp,mq}(\tau) = (1 - \eta) \int_{-\pi}^{\pi} \exp\left[-\frac{2\pi}{\lambda} (\xi'_{px} - \xi'_{qx} + \xi'_{xl} - \xi_{xm})\right] f_{\text{BS}}(x)dx + j2\pi f_D \cos(v - \gamma) f_{\text{BS}}(x)dx + \eta \int_{-\pi}^{\pi} \exp\left[-\frac{2\pi}{\lambda} (\xi_{py} - \xi_{qy} + \xi_{yl} - \xi_{ym})\right] f_{\text{MS}}(y)dy + j2\pi f_D \cos(y - \gamma) f_{\text{MS}}(y)dy \quad (4)
$$

where $\xi'$'s depend on $\phi'_k$, and $\xi$'s depend on $\phi_i$, according to Fig. 2.

This definition is the same as the one in [5] if $\tau$ is replaced by $-\tau$.  

---

Fig. 2. Detailed version of Fig. 1.
For any given \( f_{BS}(\cdot) \) and \( f_{MS}(\cdot) \), (4) can be calculated numerically using the trigonometric function relationships given in (26a)–(26h), (30a), and (30b). Note that (4) includes two parts; the first one corresponds to the STCC contributed by the scattering ring around the BS, and the second is from the scattering ring around the MS.

In order to further simplify (4), similarly to [5], the assumptions of \( D \gg \max(R', R) \), \( R' \gg \delta_{pq} \) and \( R \gg d_{lm} \) are made.\(^6\)

With these assumptions and the identities given in (26a)–(26h), (4) can be approximated by

\[
\rho_{p,mq}(\tau) \approx (1 - \eta) \int_{-\pi}^{\pi} \pi \exp \{2\pi \delta_{pq} \cos(\alpha_{pq} - x) \\
+ d_{lm} \cos(v - \beta_{lm}) \} / \lambda \\
+ \eta \int_{-\pi}^{\pi} \pi \exp \{2\pi \delta_{pq} \cos(\alpha_{pq} - w) \\
+ d_{lm} \cos(y - \beta_{lm}) \} / \lambda \\
+ \eta \int_{-\pi}^{\pi} \pi \exp \{2\pi f_{D} \cos(v - \gamma) \} \frac{f_{BS}(x)dx}{\lambda} \\
+ \eta \int_{-\pi}^{\pi} \pi \exp \{2\pi f_{D} \cos(y - \gamma) \} \frac{f_{MS}(y)dy}{\lambda}
\]

(5)

The details of the derivation are given in Appendix A.

In this paper, we consider the empirically verified von Mises angular pdf [25] for both the AoD and AoA.\(^7\) They are given by

\[
f_{BS}(x) = \frac{\exp[\kappa' \cos(x - \mu')]}{2\pi I_0(\kappa')}, \quad x \in [0, 2\pi)
\]

(6a)

\[
f_{MS}(y) = \frac{\exp[\kappa \cos(y - \mu)]}{2\pi I_0(\kappa)}, \quad y \in [0, 2\pi).
\]

(6b)

In (6a) and (6b), \( I_k(z) = (1/\pi) \int_0^\pi e^{z \cos\theta} \cos(k\theta) d\theta \) is the \( k \)-th order modified Bessel function of the first kind, \( \mu' \) and \( \mu \) account for the mean AoD and mean AoA, respectively, and \( \kappa' \) and \( \kappa \) (\( \geq 0 \)) control the angular spreads of the AoD and AoA, respectively. If \( \kappa = \kappa' = 0 \), both pdfs in (6) simplify to 1/2\( \pi \), which represents isotropic scattering. As shown in Appendix A, the mathematical form of the von Mises pdf is convenient for the analytic calculations and derivations of closed-form expressions.

Based on (30a), (30b), (33), and the assumption of \( D \gg \max(R', R) \), the following closed-form expression for the STCC is derived in Appendix A:

\[
\rho_{p,mq}(\tau) \approx (1 - \eta) \frac{\exp[-\eta(b_{lm} \cos \beta_{lm} - a \cos \gamma)]}{I_0(\kappa')} \\
\times I_0 \left\{ (\kappa^2 - a^2 \Delta^2 \sin^2 \gamma - b_{lm}^2 \Delta^2 \sin^2 \beta_{lm} \\
- c_{pq}^2 - 2b_{lm} c_{pq} \Delta' \sin \alpha_{pq} \sin \beta_{lm} \\
+ 2a \Delta' \sin \gamma (c_{pq} \sin \alpha_{pq} + b_{lm} \Delta \sin \beta_{lm}) \\
- 2b_{lm}' (a \Delta' \sin \mu' \sin \gamma - b_{lm} \Delta' \sin \beta_{lm}) \\
\times \sin \mu' - c_{pq} \cos(\alpha_{pq} - \mu')) \right\}^{1/2}
\]

(7)

Here, we have \( a = -2\pi f_{D} \tau \), \( b_{lm} = 2\pi d_{lm} / \lambda \), and \( c_{pq} = 2\pi \delta_{pq} / \mu \). Furthermore, \( 2\Delta \) is the maximum angle spread at the MS, which is determined by the scattering ring around the BS. Similarly, \( 2\Delta' \) is the maximum angle spread at the BS, which is dictated by the scattering ring around the MS. Note that (7) is a closed-form STCC function between any two subchannels of a MIMO system with arbitrary array configurations, i.e., Tx and Rx arrays do not have to be linear. In what follows, we show that many existing correlation models can be considered as special cases of (7).

1) If there is no scatterer around the BS as in a macrocell (\( \eta = 1 \)), the first half of (7) disappears, and the remaining part is the same as (12) in [5], when \( \tau \) is changed to \( -\tau \). It implies that (7) includes the model of [5] and, subsequently, other models listed in [5] as special cases.

2) With \( l = m \) and \( p = q \), the temporal autocorrelation of the subchannel \( h_{lp}(t) \) can be derived from (7) as

\[
\rho(\tau) = \frac{I_0(\sqrt{\kappa^2 - a^2 - 2\kappa \cos(\mu - \gamma)})}{I_0(\kappa)} + (1 - \eta) \\
\times I_0 \left( \frac{\sqrt{\kappa'^2 - a^2 \Delta^2 \sin^2 \gamma - 2b_{lm}' a \Delta' \sin \mu' \sin \gamma}}{e^{-\eta(b_{lm} \cos \beta_{lm} - a \cos \gamma)} I_0(\kappa') \right)
\]

(8)

where \( \rho(\tau) = \rho_{p,l,p}(\tau) \), \( \forall l, p \). If \( \eta = 1 \), (8) reduces to the model of [25].

a) With \( \eta = 1 \) and \( \kappa = 0 \) (isotropic scattering around the MS), (8) simplifies to the well-known Clarke’s temporal correlation model, i.e., \( J_0(2\pi f_{D} \tau) \) [26], where \( J_0(\cdot) \) is the Bessel function of the first kind of order zero.

\(^6\)These assumptions can be justified in many propagation scenarios of practical interest. In certain picocell propagation environments, however, \( D \gg \max(R', R) \) may not strictly hold. This may result in some discrepancies between the model and measured data in some cases, as discussed in Section V-E.

\(^7\)In this paper, no assumption is made about the antenna patterns. In fact, \( g_{l}'s \) and \( g_{l} \)'s include the Tx and Rx antenna gains, and the angular pdfs in (6a) and (6b) include the effect of antenna patterns.
3) When the MS does not move, for example, in indoor environments, one gets $f_D = 0$. This reduces (7) to the following spatial correlation between $h_{lp}$ and $h_{mq}$:

$$\rho_{lp,mq} = (1-\eta) e^{-\beta_{lm} \cos \beta_{lm}} \frac{I_0(k')}{I_0(k)} \times I_0\left(k'^2 - \frac{b_{lp}^2 \Delta^2}{\sigma_{pq}^2} \sin^2 \beta_{lm} - \frac{2b_{lm} c_{pq} \Delta}{\sigma_{pq}} \sin \alpha_{pq} \sin \beta_{lm} + j2\kappa' \sin \alpha_{pq} \sin \beta_{lm} + j2\kappa' \sin \alpha_{pq} \sin \beta_{lm} + j2\kappa' \cos(\alpha_{pq} - \mu') \right) + \eta e^{j\pi \rho_{pq} \cos \alpha_{pq}} + \frac{1}{I_0(k)} \times I_0\left(k'^2 - \frac{b_{lp}^2 \Delta^2}{\sigma_{pq}^2} \sin^2 \beta_{lm} - \frac{2b_{lm} c_{pq} \Delta}{\sigma_{pq}} \sin \alpha_{pq} \sin \beta_{lm} + j2\kappa' \sin \alpha_{pq} \sin \beta_{lm} + j2\kappa' \sin \alpha_{pq} \sin \beta_{lm} + j2\kappa' \cos(\alpha_{pq} - \mu') \right) \right)^2. \tag{9}$$

With $p = q$ and $\eta = 1$, (10) reduces to $\rho_{lp,mp} = J_0(2\pi d_{lm}/\lambda)$. It is the spatial correlation at the MS in a macrocell [26]. On the other hand, with $l = m$ and $\eta = 1$, (10) simplifies to $\rho_{lp,lq} = J_0(\Delta 2\pi \delta_{pq}/\lambda)$. It is the spatial correlation at the BS in a macrocell [27].

IV. DATA PROCESSING AND PARAMETER ESTIMATION

This section explains how the raw data are processed and how the parameters of the proposed model are estimated.

A. Normalization

First, it is necessary to do some power normalization. In [28], the normalization is done according to $\tilde{h}(t) = \sqrt{N_T N_R} H(t)/\|H(t)\|_F$ for each $t$, $t \in [1, T]$, where $H(t)$ is a snapshot of the channel matrix at time instant $t$, $T$ is the total number of snapshots, and $\|H(t)\|_F = \sqrt{\sum_{i=1}^{N_T} \sum_{p=1}^{N_R} |h_{lp}(t)|^2}$. Note that $H(t)$ denotes $H(t_{T_{x}}, t_{T_{s}})$ and $T_{s}$ is the sampling period. To simplify notations, $T_{s}$ is dropped in this section. In [29], normalization is intended to guarantee the so-called unit single-input single-output gain, i.e., $T^{-1} \sum_{t=1}^{T} \|H(t)\|_F^2 = N_T N_R$; therefore, the normalization is done according to

$$\tilde{h}(t) = \sqrt{\frac{N_R N_T T}{\sum_{t=1}^{T} \|H(t)\|_F^2}} H(t), \quad t \in [1, T]. \tag{11}$$

According to our correlation model in (7), each individual subchannel has unit power. On the other hand, the aforementioned two normalization methods do not result in unit power for each subchannel. Hence, in this paper, normalization is performed such that each subchannel has unit variance and zero mean, i.e.,

$$\tilde{h}_{lp}(t) = \frac{h_{lp}(t) - \hat{m}_{lp}}{\sigma_{lp}}, \quad \forall l, p, t \tag{12}$$

where $\hat{m}_{lp}$ and $\sigma_{lp}^2$ are the estimated mean and variance of the BS-Ms subchannel. They are calculated from $\{h_{lp}(t)\}_{t=1}^{T}$. According to $\hat{m}_{lp} = T^{-1} \sum_{t=1}^{T} h_{lp}(t)$ and $\sigma_{lp} = T^{-1} \sum_{t=1}^{T} |h_{lp}(t) - \hat{m}_{lp}|^2$.

B. Parameter Estimation

First, we need to determine $\alpha$ and $\beta$, which are the directions of BS and MS arrays, respectively. The direction of the MS motion $\gamma$ and the Doppler drift $f_D$ in the HUT data needs to be obtained as well, whereas the fixed MS in the BYU data does not need $\gamma$ and $f_D$.

1) For the HUT data, a comparison of Figs. 12 and 13 with Fig. 1 gives $\alpha = \pi/2, \beta = 4\pi/5$, and $\gamma = \pi/2$. To obtain these angles, first, $x$- and $y$-axes were chosen in the horizontal and vertical directions, respectively, similar to Fig. 1, and were redrawn on Fig. 12. This results in Fig. 3. A comparison of Fig. 3 with Fig. 1 reveals that the BS array is parallel to the $y$-axis; therefore, $\alpha = \pi/2$. $\beta$ can be determined similarly as $144^\circ = 180^\circ - (90 - 54)^\circ$ based on the $54^\circ$ angle shown in Fig. 3. A further comparison of Fig. 3 with Fig. 1 shows that the mobile is moving parallel to the $y$-axis in the positive direction. Therefore, $\gamma = \pi/2$. The data are taken over a 47-m route in Unioninkatu Street, at the intersection with Rauhankatu Street, as shown in Fig. 12. $f_D = 2.872$ Hz, which is derived in Appendix B. Note that only a small section

\[\]
To estimate the rest of the parameters $\Lambda = (\Delta, \Delta', \kappa, \kappa', \mu, \mu', \eta)$ of the model, a nonlinear least-square correlation fitting approach is used, via a numerical search over the parameter space. This is explained next.

For the BYU data, $\Lambda$ is estimated\(^\text{10}\) by fitting the spatial MIMO correlation matrix $R(\Lambda)$, whose elements are given by (9), to the estimated spatial correlation matrix $\hat{R}$

$$
\hat{\Lambda} = \arg \min_{\Lambda} \left\| \hat{R} - R(\Lambda) \right\|_F^2.
$$

\(^{10}\)The function in (13) is very complex since it depends on the seven free variables. The same observation applies to (14). Finding the optimal estimates is not the focus of this paper. This is a separate issue which is recently studied in [30].

Similarly, for the HUT data, $\Lambda$ is estimated by fitting the spatio-temporal MIMO correlation matrix $R(\Lambda, k)$, whose elements are determined by (7), to the estimated spatio-temporal correlation matrix $\hat{R}(k)$

$$
\hat{\Lambda} = \arg \min_{\Lambda} \sum_{k=0}^{t_{\text{max}}} \left\| \hat{R}(k) - R(\Lambda, k) \right\|_F^2
$$

where $t_{\text{max}}$ is set to 4 in the numerical search.\(^\text{11}\) $\hat{R}(k) = T^{-1} \sum_{t=1}^T h(t+k)h^\dagger(t)$, $k \in [0, t_{\text{max}}]$, $R(\Lambda, k) = \mathbb{E}[h(t+k)h^\dagger(t)]$ such that $R(\Lambda, k)_{l+N_k(p-1),m+N_k(q-1)} = \rho_{lp,mq}$, $l, m \in [1, N_R]$, and $p, q \in [1, N_T]$. Note that the expectation is taken over AoA and AoD, as shown in (4), and $\rho_{lp,mq}$ is given by (9). Furthermore, $R(\Lambda)$ is an $N_T N_R \times N_T N_R$ Hermitian block Toeplitz with Toeplitz blocks (H-BTTB) matrix. It has $N_T \times N_T$ blocks, and each block is an $N_R \times N_R$ square matrix. A matrix is BTTB if its $(i, j)$th block is a function of $(i-j)$, and the $(i,j)$th block itself is a Toeplitz matrix. More discussion on BTTB matrices can be found in [31]. The Hermitian property is due to the Hermitian symmetry of the STCC $\rho_{lp,mq}(\tau)$ around $\tau = 0$. The BTTB structure comes from the fact that $\rho_{lp,mq}(\tau) = \rho((i+j)+(p+j)-(m+i)\tau)$, $\forall i, j$. By taking advantage of the H-BTTB structure of $R(\Lambda)$, the numerical search in (13) can be performed much faster, particularly for large $N_R$ and $N_T$.

Similarly, for the HUT data, $\Lambda$ is estimated by fitting the spatio-temporal MIMO correlation matrix $R(\Lambda, k)$, whose elements are determined by (7), to the estimated spatio-temporal correlation matrix $\hat{R}(k)$

$$
\hat{\Lambda} = \arg \min_{\Lambda} \sum_{k=0}^{t_{\text{max}}} \left\| \hat{R}(k) - R(\Lambda, k) \right\|_F^2
$$

where $t_{\text{max}}$ is set to 4 in the numerical search.\(^\text{11}\) $\hat{R}(k) = T^{-1} \sum_{t=1}^T h(t+k)h^\dagger(t)$, $k \in [0, t_{\text{max}}]$, $R(\Lambda, k) = \mathbb{E}[h(t+k)h^\dagger(t)]$ such that $R(\Lambda, k)_{l+N_k(p-1),m+N_k(q-1)} = \rho_{lp,mq}(k)$, $l, m \in [1, N_R]$, and $p, q \in [1, N_T]$. Note that $\rho_{lp,mq}(k)$ denotes $\rho_{lp,mq}(kT_s)$ given by (7). In addition, it is clearly seen that $R(\Lambda, 0)$ is an H-BTTB matrix, whereas for $k \neq 0$, $R(\Lambda, k)$ is a BTTB matrix. In a similar manner to (13), the parameter search in (14) can be sped up using the BTTB structure of $R(\Lambda, k)$.

It is clear from (9) that if $\mu$ and $\mu'$ are the outcomes of the search in (13), $2\beta_{lm} - \hat{\mu}$ and $2\alpha_{pq} - \hat{\mu}'$ are the valid parameters

\(^{11}\)According to Appendix B, $t_{\text{max}} = 4$ corresponds to $1/f_p$ sec.

as well and result in the same numerical values for (9). On the other hand, if $\hat{\mu}'$ is the search result of (14), $2\alpha_{pq} - \hat{\mu}'$ gives the same numerical value for (7) as well. All these angular ambiguities should be considered with the real propagation environments when looking at the search results of (13) and (14). Moreover, for the HUT data, the nonuniform antenna patterns [22, Fig. 2] were also considered when searching for parameters.\(^\text{7}\) Finally, the estimated parameter sets for the HUT and BYU data are presented in Table I.

### V. COMPARISON OF THE PROPOSED MODEL WITH COLLECTED DATA

This section presents the comparison of the proposed model with collected data in terms of various aspects of fading channels such as the statistical distribution of amplitude, phase, in-phase/quadrature components, spatial, temporal, and spatio-temporal correlations, LCR, AFD, and MIMO mutual information.

#### A. Statistical Distribution of the Data

When deriving the compact STCC expression in (7), it has been assumed that the numbers of local scatterers around both the BS and MS are large enough. This translates into complex Gaussian distribution for each subchannel $h_{lp}(t)$ according to the central limit theorem [32]. To verify this, the data are normalized according to (12) such that each subchannel has zero mean with unit variance. As shown in Fig. 4, the real and imaginary parts of $h_{11}(t)$ of the HUT data are Gaussian. Furthermore, the empirical cumulative distribution functions (CDFs) of the amplitude and phase are Rayleigh and uniform, respectively. The same types of plots are observed for other HUT subchannels, as well as the BYU data sets, as reported in [33, p. 162] and [34, Fig. 5].

#### B. Spatial Correlations

Here, we consider four different types of spatial correlations, i.e., parallel, crossing, transmit, and receive correlations, which are defined by

$$
\rho_{\text{parallel}} \overset{\Delta}{=} \mathbb{E} \left[ h_{lp}(t)h^\dagger_{lp}(t+u)(t) \right] \quad (15a)
$$

$$
\rho_{\text{crossing}} \overset{\Delta}{=} \mathbb{E} \left[ h_{lp}(t)h^\dagger_{mp}(t) \right] \quad (15b)
$$

$$
\rho_{\text{Tx}} \overset{\Delta}{=} \mathbb{E} \left[ h_{lp}(t)h^\dagger_{mp}(t) \right] \quad (15c)
$$

$$
\rho_{\text{Rx}} \overset{\Delta}{=} \mathbb{E} \left[ h_{lp}(t)h^\dagger_{mp}(t) \right] \quad (15d)
$$

\begin{table}[h]
\centering
\caption{Estimated Parameter Sets for the HUT and BYU Data}
\begin{tabular}{c c c c c c c c}
\hline
 & $\hat{\Delta}$ & $\hat{\Delta}'$ & $\hat{\kappa}$ & $\hat{\kappa}'$ & $\hat{\mu}$ & $\hat{\mu}'$ & $\hat{\eta}$ \\
\hline
HUT & $\frac{\pi}{6}$ & $\frac{\pi}{4}$ & 17 & 2 & $\frac{15\pi}{2}$ & 0.7 \\
BYU & $\frac{\pi}{3}$ & $\frac{\pi}{6}$ & 0 & 0.5 & $\frac{5\pi}{8}$ & 0.2 \\
\hline
\end{tabular}
\end{table}
In a $2 \times 2$ channel, it can be shown that
\begin{align}
\rho_{\text{parallel}} & \triangleq \mathbb{E} [h_{11}(t)h_{22}^*(t)] = \rho_{11,22} \quad (16) \\
\rho_{\text{crossing}} & \triangleq \mathbb{E} [h_{12}(t)h_{21}^*(t)] = \rho_{12,21} \quad (17) \\
\rho_{\text{Tx}} & \triangleq \mathbb{E} [h_{11}(t)h_{12}^*(t)] = \mathbb{E} [h_{21}(t)h_{22}^*(t)] = \rho_{11,12} = \rho_{21,22} \quad (18) \\
\rho_{\text{Rx}} & \triangleq \mathbb{E} [h_{11}(t)h_{21}^*(t)] = \mathbb{E} [h_{12}(t)h_{22}^*(t)] = \rho_{11,21} = \rho_{12,22}. \quad (19)
\end{align}

For the HUT data, the model correlations are obtained as
$(\rho_{\text{parallel}}, \rho_{\text{crossing}}, \rho_{\text{Tx}}, \rho_{\text{Rx}}) = (0.04, 0.1, 0.6, 0.08)$
by plugging $\hat{\Lambda}_{\text{HUT}}$ from Table I into (9), whereas the empirical correlations are
$(0.04, 0.1, 0.6, 0.08)$ by processing the normalized data. For the BYU data, the same type of comparison for
the aforementioned four correlations in terms of element spacing is shown in Fig. 5. The utility of the model is
revealed by the close agreement between the model and the empirical spatial correlations.

C. Spatio-Temporal Cross Correlations

Fig. 6 shows the empirical autocorrelation of the subchannel $h_{11}(t)$ for the HUT data, together with the autocorrelation of the model obtained by plugging $\hat{\Lambda}_{\text{HUT}}$ of Table I into (8). The STCCs of the HUT data between $h_{11}(t)$ and $h_{12}(t)$ and
between $h_{12}(t)$ and $h_{11}(t)$, respectively, are shown in Figs. 7 and 8, where the cross correlation of the model is plotted by
substituting $\hat{\Lambda}_{\text{HUT}}$ of Table I into (7). The difference between the model cross correlations in Figs. 7 and 8 is due to
$\alpha_{21} = \alpha_{12} + \pi$, as explained in footnote 3. Clearly, the model provides a close fit to the empirical spatio-temporal correlations.

D. Level Crossing Rate and Average Fade Duration

LCR and AFD of the signal envelope are two important temporal statistical features which provide useful
information on the dynamic behavior of time-varying fading channels. To calculate the LCR and AFD of a subchannel, one needs its
temporal autocorrelation. This is given in (8). With $\gamma = \pi/2$
in the HUT measurement campaign, (8) reduces to
\begin{equation}
\rho(\tau) = (1 - \eta) \frac{I_0 \left( \sqrt{\kappa^2 - a^2 \Delta^2 - 2 \kappa \alpha \Delta \sin \mu} \right)}{I_0(\kappa')} + \eta \frac{I_0 \left( \sqrt{\kappa^2 - a^2 - 2 \kappa \alpha \sin \mu} \right)}{I_0(\kappa)}. \quad (20)
\end{equation}

The $k$th spectral moment $B_k$ is defined by [35]
\begin{equation}
B_k = \frac{d^k \rho(\tau)}{d\tau^k} \bigg|_{\tau=0} \quad (21)
\end{equation}
which is also required for calculating LCR and AFD. From (20) and (21), we obtain
\begin{align}
B_0 &= 1 \quad (22a) \\
B_1 &= 2\pi f_D \left[\frac{(1 - \eta)I_1(\kappa') \Delta' \sin \mu'}{I_0(\kappa')} + \frac{\eta I_1(\kappa) \sin \mu}{I_0(\kappa)}\right] \quad (22b) \\
B_2 &= 4\pi^2 f_D^2 \left\{ \eta \left[ \sin^2 \mu + \frac{I_1(\kappa') \cos (2\mu')}{\kappa I_0(\kappa')} \right] + (1 - \eta) \Delta'^2 \left[ \sin^2 \mu' + \frac{I_1(\kappa') \cos (2\mu')}{\kappa' I_0(\kappa')} \right] \right\}. \quad (22c)
\end{align}
Fig. 5. Comparison of four different types of spatial correlations of the new MIMO model with the corresponding empirical correlations of the BYU data collected on November 7, 2000.

Fig. 6. Autocorrelation of $h_{11}(t)$, which is defined by $E[h_{11}(t+\tau)h_{11}^*(t)]$. New model and HUT data.

The theoretical LCR and AFD at the given threshold $r$ for Rayleigh fading are, respectively, given by [35]

$$N_{|h|}(r) = \sqrt{\frac{B_2 - B_1^2}{\pi}} r e^{-r^2}$$  \hspace{1cm} (23)$$

$$\tilde{t}_{|h|}(r) = \frac{e^{r^2} - 1}{r \sqrt{B_2 - B_1^2}}$$  \hspace{1cm} (24)$$

where $|h|$ denotes the amplitude of any subchannel.
By plugging the estimated parameter set $\hat{\Lambda}_{HUT}$ of Table I into (23) and (24), the LCR and AFD of the model are compared in Fig. 9 with the empirical LCR and AFD of $|h_{22}(t)|$ versus $r$. The close match for this and three other subchannels, which is not shown due to space limitation, verifies the ability of the model in capturing the dynamics of the channel.

### E. CDF of Instantaneous Mutual Information (IMI)

In this section, we focus on the IMI under equal-power transmission to verify our proposed model. Assuming that the channel matrix $H(t)$ is known at the receiver but not at the transmitter, with the MIMO channel characterized in (1), the IMI under equal-power transmission, in bits per second per hertz (bps/Hz), is given by [4]

$$I(t) = \log_2 \det \left( I_{NR} + \frac{P_{Tx}}{N_T P_{Noise}} H(t) H^\dagger(t) \right)$$

(25)

where $\det(\cdot)$ is the matrix determinant. Note that at any time instant $t$, $I(t)$ is a random variable, and its distribution depends
Fig. 9. LCR and AFD of $h_{22}(t)$. New model and HUT data.

Fig. 10. Distributions of the IMI: HUT data.

on the distribution of the random matrix $H(t)H^\dagger(t)$. In this paper, we compare the empirical distribution of the MIMO IMI of the data, with the distribution predicted by our proposed model, where the SNR $= P_{Tx}/P_{Noise}$ is set to 20 dB in (25).

In Figs. 10 and 11, three different CDFs of IMI are plotted for the HUT and BYU data, respectively. The “empirical” IMI is calculated using (25) directly from the collected snapshots $\{H(t)\}_{t=1}^T$, normalized according to (12). However, for each of the other IMIs, first, a zero-mean $N_R \times N_T$ Gaussian matrix $H_{IID}$ with i.i.d. elements is generated for each $t$ such that $h_{IID} = \text{vec}(H_{IID})$ and $E[h_{IID}h_{IID}^\dagger] = I_{N_RN_T}$. The “IID” IMI CDF is obtained by inserting $H_{IID}$ into (25). To obtain the “new model” IMI CDF, the following necessary steps are done.

1) For the HUT data, the correlated matrix-valued random process $H(t)$ is generated using the spectral method [36] based on (7) and $\hat{\Lambda}_{HUT}$ of Table I and then plugged into (25).

2) For the BYU data, the colored Gaussian vector $h(\hat{\Lambda}_{BYU})$ is built according to $h(\hat{\Lambda}_{BYU}) = \sqrt{R(\hat{\Lambda}_{BYU})}h_{IID}$
where $\sqrt{\mathbf{R}(\hat{\mathbf{A}}_{\text{BYU}})}$ is the Cholesky matrix of $\mathbf{R}(\hat{\mathbf{A}}_{\text{BYU}})$, $\mathbf{R}(\mathbf{A})$ is defined in Section IV-B, and $\hat{\mathbf{A}}_{\text{BYU}}$ is given in Table I. Then, the channel matrix $\mathbf{H}(\hat{\mathbf{A}}_{\text{BYU}})$ is generated from $\mathbf{h}(\hat{\mathbf{A}}_{\text{BYU}}) = \text{vec}(\mathbf{H}(\hat{\mathbf{A}}_{\text{BYU}}))$ and (25).

According to Figs. 10 and 11, the model fits the HUT data well. However, the match to the BYU data is not as good as the HUT data. This may be partly because the assumption $D \gg \max(R', R)$ does not strictly hold for the scenarios where the BYU data were collected.

VI. Conclusion

In this paper, a parametric statistical MIMO model is proposed where both the BS and the MS experience local scattering. The model yields a closed-form and mathematically tractable expression for the STCC between any two subchannels of a MIMO system with arbitrary transmit and receive array configurations. The model includes some key physical channel parameters such as the mean AoD and AoA, anglespreads, maximum Doppler frequency, sizes of the local scattering rings, and the relative contribution of BS and MS local scatterers.

The proposed model is general enough to include, as special cases, most existing correlation models such as the MIMO model of [5], Lee’s spatio-temporal correlation model, and Clarke’s classical spatial/temporal correlation model. The comparison of the model with outdoor and indoor collected data, in terms of spatial correlations, temporal autocorrelations, STCCs, LCRs, AFD, and the distribution of mutual information, has demonstrated the flexibility of the model in describing real-world propagation environments. The proposed model, which is supported by empirical observation, provides a useful channel characterization required for the efficient design and performance prediction of multiantenna transceivers.

APPENDIX A

DERIVATION OF (5) AND (7)

Based on the application of the law of cosines in appropriate triangles in Fig. 2, we obtain

$$\xi_{px}' = \delta_{dq}/4 + R'^2 - \delta_{pq} R' \cos(\alpha_{pq} - x) \quad (26a)$$
$$\xi_{qx}' = \delta_{dq}/4 + R'^2 + \delta_{pq} R' \cos(\alpha_{pq} - x) \quad (26b)$$
$$\xi_{zl}' = d_{lm}/4 + \xi_{z}' - \xi_{z} d_{lm} \cos(v - \beta_{lm}) \quad (26c)$$
$$\xi_{xm}' = d_{lm}/4 + \xi_{x}' + \xi_{x} d_{lm} \cos(v - \beta_{lm}) \quad (26d)$$
$$\xi_{py}' = \delta_{pq}/4 + \xi_{y}' - \delta_{pq} \cos(\alpha_{pq} - w) \quad (26e)$$
$$\xi_{qy}' = \delta_{pq}/4 + \xi_{y}' + \delta_{pq} \cos(\alpha_{pq} - w) \quad (26f)$$
$$\xi_{yl}' = d_{lm}/4 + R'^2 - d_{lm} R \cos(y - \beta_{lm}) \quad (26g)$$
$$\xi_{ym}' = d_{lm}/4 + R'^2 + d_{lm} R \cos(y - \beta_{lm}) \quad (26h)$$

where, to simplify the notation, we use $x$ for $\phi_{k}'$, $y$ for $\phi_{l}'$, $v$ for $\varphi_{k}'$, $w$ for $\varphi_{l}'$, $\xi_{xi}'$ for $\xi_{xi}'$, and $\xi_{yi}'$ for $\xi_{yi}'$.

The following identity can be obtained by subtracting (26b) from (26a):

$$\xi_{px}' - \xi_{qx}' = \frac{-2 \delta_{pq} R' \cos(\alpha_{pq} - x)}{\xi_{px}' + \xi_{qx}'}. \quad (27)$$

With the assumption of $R' \gg \delta_{pq}$, we have $\xi_{px}' + \xi_{qx}' \approx 2R'$; therefore, (27) reduces to

$$\xi_{px}' - \xi_{qx}' \approx -\delta_{pq} \cos(\alpha_{pq} - x). \quad (28)$$

By applying the same reasoning to (26c)–(26h) and using the assumptions $D \gg \max(R', R)$ and $R \gg d_{lm}$, which imply
\(\xi'_{x} + \xi'_{x,m} \approx 2\xi'_{x},\ \xi_{py} + \xi_{qy} \approx 2\xi_{y},\ \text{and} \ \xi_{yl} + \xi_{ym} \approx 2R,\) we obtain the following approximations:

\[
\begin{align*}
\xi'_{x} - \xi'_{x,m} & \approx -d_{lm} \cos (v - \beta_{lm}) \\
\xi_{py} - \xi_{qy} & \approx -\delta_{pq} \cos (\alpha_{pq} - w) \\
\xi_{yl} - \xi_{ym} & \approx -d_{lm} \cos (y - \beta_{lm}).
\end{align*}
\]

Substitution of (28) and (29) into (4) results in the integral representation (5).

Now, we apply the law of sines to the triangles \(O'S'_{l}O\) and \(O'S_{l}O\), respectively, to obtain the following identities:

\[
\begin{align*}
\frac{D}{\sin(v-x)} & = \frac{R'}{\sin(\pi - v)} \quad (30a) \\
\frac{D}{\sin(y-w)} & = \frac{R}{\sin w}. \quad (30b)
\end{align*}
\]

From Fig. 2, it is clear that \(\pi - v \leq \Delta', \ \omega \leq \Delta, \ \Delta' = \arcsin(R'/D), \ \text{and} \ \Delta = \arcsin(R/D).\) Based on the assumption \(D \gg \max(R' , R),\) we conclude that \(\Delta'\) and \(\Delta\) and, consequently, \(\pi - v\) and \(\omega\) are small quantities. This observation, together with \(\sin \epsilon \approx \epsilon\) when \(\epsilon\) is small, allows us to derive the following approximations from (30a) and (30b):

\[
\begin{align*}
v & \approx \pi - \Delta' \sin x \\
\omega & \approx \Delta \sin y.
\end{align*}
\]

Furthermore, using \(\sin \epsilon \approx \epsilon\) and \(\cos \epsilon \approx 1\) when \(\epsilon\) is small, together with (31a) and (31b), \(\cos (v - \beta_{lm})\) and \(\cos (\alpha_{pq} - w)\) can be approximated as

\[
\begin{align*}
\cos (v - \beta_{lm}) & \approx -\cos \beta_{lm} + \Delta' \sin \beta_{lm} \sin x \\
\cos (\alpha_{pq} - w) & \approx \cos \alpha_{pq} + \Delta \sin \alpha_{pq} \sin y.
\end{align*}
\]

By plugging (6b), (6a), (32a), and (32b) into (5) and calculating the two integrals according to [37, 3.38-4]

\[
\int_{-\pi}^{\pi} \exp(x \sin \theta + y \cos \theta) d\theta = 2\pi I_{0} \left(\sqrt{x^{2} + y^{2}}\right)
\]

the general STCC in (7) can be obtained after some algebraic manipulations.

**APPENDIX B**

**HUT DATA**

Fig. 12 shows the layout of the location where the HUT data are collected. The \(2 \times 2\) measurement setup and four subchannels are shown in Fig. 13(a), whereas the configuration of the Rx array is presented in Fig. 13(b). The Tx and Rx element spacings are given by \(d_{12} = \lambda\) and \(d_{12} = 2.785\lambda\), respectively. The mobile speed is 0.4 m/s, which, with the 2.154-GHz carrier frequency, results in \(f_{D} = \nu/\lambda = 2.872\) Hz. There are \(T = 1342\) snapshots of the channel matrix \(H(t)\), and any two adjacent snapshots are separated by \(\lambda/4\) in space, which is equivalent to 87.05 ms in time. Further details of the data and measurement setup can be found in [21] and [22]. The pathloss is removed from the data by least-square fitting of a tenth order polynomial in the decibel domain, and the power normalization is performed according to (12).

**APPENDIX C**

**BYU DATA**

For the data collected on November 7, 2000, at BYU, the room layout is shown in Fig. 14. Both the Tx and Rx arrays are linear, each with ten monopole elements and \(\lambda/4\) spacing at the carrier frequency of 2.42 GHz. The Tx array was placed at four different locations in room 484, whereas the Rx array was placed at six different locations in room 400. There are 24 Tx–Rx location configurations, 20 data sets per location, and 124 snapshots of the \(10 \times 10\) channel matrix per data set. Power normalization of each data set was done according to

\[
\tilde{h}_{loc, set, lp}(t) = \frac{h_{loc, set, lp}(t) - \mu_{loc, set, lp}}{\sigma_{loc, set, lp}}
\]

where \(loc\) is the location index \([1, 20]\), \(l, p \in [1, 10]\), \(t \in [1, 124]\).
loc\textsuperscript{th} location. Because of the fixed Tx and Rx and stationary environment, we have $f_D = 0$ Hz. More measurement information can be found in [33, ch. 6 and 7] and [34], where this data set is labeled as “10 × 10(a).”

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